

**IN THE UNITED STATES PATENT AND TRADEMARK OFFICE**

In re Patent of: Jeroen Poeze et al.  
U.S. Patent No.: 10,631,765 Attorney Docket No.: 50095-0024IP1  
Issue Date: April 28, 2020  
Appl. Serial No.: 16/725,478  
Filing Date: December 23 2019  
Title: MULTI-STREAM DATA COLLECTION SYSTEM FOR  
NONINVASIVE MEASUREMENT OF BLOOD  
CONSTITUENTS

**THIRD DECLARATION OF JACOB ROBERT MUNFORD**

1. My name is Jacob Robert Munford. I am over the age of 18, have personal knowledge of the facts set forth herein, and am competent to testify to the same.
  
2. I earned a Master of Library and Information Science (MLIS) from the University of Wisconsin-Milwaukee in 2009. I have over ten years of experience in the library/information science field. Beginning in 2004, I have served in various positions in the public library sector including Assistant Librarian, Youth Services Librarian and Library Director. I have attached my Curriculum Vitae as Appendix CV.
  
3. During my career in the library profession, I have been responsible for materials acquisition for multiple libraries. In that position, I have cataloged, purchased and processed incoming library works. That includes purchasing materials directly from vendors, recording publishing data from the material in question, creating detailed material records for library catalogs and physically preparing that material for circulation. In addition to my experience in acquisitions, I was also responsible for analyzing large collections of library materials, tailoring library records for optimal catalog

search performance and creating lending agreements between libraries during my time as a Library Director.

4. I am fully familiar with the catalog record creation process in the library sector. In preparing a material for public availability, a library catalog record describing that material would be created. These records are typically written in Machine Readable Catalog (herein referred to as “MARC”) code and contain information such as a physical description of the material, metadata from the material’s publisher, and date of library acquisition. In particular, the 008 field of the MARC record is reserved for denoting the date of creation of the library record itself. As this typically occurs during the process of preparing materials for public access, it is my experience that an item’s MARC record indicates the date of an item’s public availability.
  
5. Typically, in creating a MARC record, a librarian would gather various bits of metadata such as book title, publisher and subject headings among others and assign each value to a relevant numerical field. For example, a book’s physical description is tracked in field 300 while title/attribution is tracked in field 245. The 008 field of the MARC record is reserved for denoting the creation of the library record itself. As this is the only date reflecting the inclusion of said materials within the library’s collection, it is my experience

that an item's 008 field accurately indicates the date of an item's public availability.

6. This declaration is being drafted as of September 2021. Public and university libraries in my area have been closed for months due to the COVID-19 pandemic. My state, Pennsylvania, has a travel advisory, which has affected my ability to travel. In my experience, library catalog records are accurate descriptions of a library's collection and my lack of physical access to libraries at this time creates no doubt in my determinations of authenticity or availability of the exhibits noted below.
  
7. I have reviewed Exhibit APPLE-1055, "Refractive Indices of Human Skin Tissues at Eight Wavelengths and Estimated Dispersion Relations Between 300 and 1600 nm" by Huafeng Ding, Jun Lu, William Wooden, Peter Kragel and Xin-Hua Hu as published in *Physics in Medicine & Biology* in March of 2006.
  
8. Attached hereto as Appendix DING01, is a true and correct copy of scans from *Physics in Medicine & Biology* as held by the Penn State University library. I secured these scans myself from the library's collection. In comparing Exhibit APPLE-1055 to Appendix DING01, it is my



determination that Exhibit APPLE-1055 is a true and correct copy of “Refractive Indices of Human Skin Tissues at Eight Wavelengths and Estimated Dispersion Relations Between 300 and 1600 nm” by Huafeng Ding, Jun Lu, William Wooden, Peter Kragel and Xin-Hua Hu as published in *Physics in Medicine & Biology* in March of 2006.

9. Attached hereto as Appendix DING02 is a true and correct copy of the MARC record for *Physics in Medicine & Biology* as held by the Penn State University library. I secured this record myself from the library’s public catalog. The MARC record contained within Appendix DING02 accurately describes the title, author, publisher, and ISSN number of *Physics in Medicine & Biology*. The ‘Availability’ field of this record visible on pages 7-11 indicates this collection includes the March 2006 edition of *Physics in Medicine & Biology* containing “Refractive Indices of Human Skin Tissues at Eight Wavelengths and Estimated Dispersion Relations Between 300 and 1600 nm”.

10. The 008 field of each MARC record in Appendices DING02 indicates the date of record creation. The 008 field of Appendix DING02 indicates the Penn State University library first cataloged this journal as of August 5, 1982 and held the journal in perpetuity until 2012. Considering this

information, it is my determination that *Physics in Medicine & Biology* March 2006 and therefore “Refractive Indices of Human Skin Tissues at Eight Wavelengths and Estimated Dispersion Relations Between 300 and 1600 nm” was made available to the public shortly after its initial release in March of 2006.

11. I have reviewed Exhibit APPLE-1051, *Optics* by Eugene Hecht, 2<sup>nd</sup> Edition.

12. Attached hereto as Appendix HECHT01, is a true and correct copy of scans of select pages from *Optics, 2<sup>nd</sup> Edition* textbook as held by the University of Pittsburgh library. I secured the textbook and made these scans myself from the library’s collection. For the purposes of this declaration, these scans do not show the entire textbook, but only select pages of the textbook. In comparing Exhibit APPLE-1051 to the textbook from which I obtained the partial scans of Appendix HECHT01, it is my determination that Exhibit APPLE-1051 is a true and correct copy of *Optics, 2<sup>nd</sup> Edition*.

13. Attached hereto as Appendix HECHT02 is a true and correct copy of the MARC record for *Optics, 2<sup>nd</sup> Edition* as held by the University of Pittsburgh library. I secured this record myself from the library’s public catalog. The

MARC record contained within Appendix HECHT02 accurately describes the title, author, publisher, and ISBN number of *Optics, 2<sup>nd</sup> Edition*.

14. The 008 field of each MARC record in Appendices HECHT02 indicates the date of record creation. The 008 field of Appendix HECHT02 indicates the University of Pittsburgh library first acquired this book as of March 3, 1994. Considering this information, it is my determination that *Optics, 2<sup>nd</sup> Edition* was made available to the public shortly after its initial release in Spring of 1994.

15. I have reviewed Exhibit APPLE-1049, *Optics* by Eugene Hecht, 4<sup>th</sup> Edition.

16. Attached hereto as Appendix HECHT03, is a true and correct copy of scans of select pages from *Optics, 4<sup>th</sup> Edition* textbook as held by the University of Pittsburgh library. I secured the textbook and made these scans myself from the library's collection. For the purposes of this declaration, these scans do not show the entire textbook, but only select pages of the textbook. In comparing Exhibit APPLE-1049 to the textbook from which I obtained the partial scans of Appendix HECHT03, it is my determination that Exhibit APPLE-1049 is a true and correct copy of *Optics, 4<sup>th</sup> Edition*.

17. Attached hereto as Appendix HECHT04 is a true and correct copy of the MARC record for *Optics, 4<sup>th</sup> Edition* as held by the University of Pittsburgh library. I secured this record myself from the library's public catalog. The MARC record contained within Appendix HECHT04 accurately describes the title, author, publisher, and ISBN number of *Optics, 4<sup>th</sup> Edition*.

18. The 008 field of each MARC record in Appendices HECHT04 indicates the date of record creation. The 008 field of Appendix HECHT04 indicates the University of Pittsburgh library first acquired this book as of November 5, 2001. Considering this information, it is my determination that *Optics, 4<sup>th</sup> Edition* was made available to the public shortly after its initial release in Fall of 2001.

19. I have reviewed Exhibit APPLE-1056, "Analysis of the Dispersion of Optical Plastic Materials" by Stefka Kasarova, Nina Sultanova, Christo Ivanov and Ivan Nikolov as published in *Optical Materials* in July of 2007.

20. Attached hereto as Appendix KASAROVA01, is a true and correct digital copy of "Analysis of the Dispersion of Optical Plastic Materials" by Stefka Kasarova, Nina Sultanova, Christo Ivanov and Ivan Nikolov as held by the University of Notre Dame library. I secured this copy myself from the

library's collection. In comparing Exhibit APPLE-1056 to Appendix KASAROVA01, it is my determination that Exhibit APPLE-1056 is a true and correct copy of "Analysis of the Dispersion of Optical Plastic Materials" by Stefka Kasarova, Nina Sultanova, Christo Ivanov and Ivan Nikolov as published in *Optical Materials* in July of 2007.

21. Attached hereto as Appendix KASAROVA02 is a true and correct copy of the MARC record for *Optical Materials* as held by the University of Notre Dame library. I secured this record myself from the library's public catalog. The MARC record contained within Appendix KASAROVA02 accurately describes the title, author, publisher, and ISSN number of *Optical Materials*. Selections from the table of contents of this record visible on page 4 indicates this collection includes the July 2007 edition of *Optical Materials* containing "Analysis of the Dispersion of Optical Plastic Materials". In comparing Exhibit APPLE-1056 to Appendix KASAROVA02, it is my determination that Exhibit APPLE-1056 is a true and correct copy of "Analysis of the Dispersion of Optical Plastic Materials" by Stefka Kasarova, Nina Sultanova, Christo Ivanov and Ivan Nikolov, as published in *Optical Materials* in July of 2007.

22. The 008 field of each MARC record in Appendices KASAROVA02 indicates the date of record creation. The 008 field of Appendix KASAROVA02 indicates the University of Notre Dame library first cataloged this journal as of May 11, 1998 and still holds the journal in perpetuity. Considering this information, it is my determination that *Optical Materials* July 2007 and therefore “Analysis of the Dispersion of Optical Plastic Materials” was made available to the public shortly after its initial release in July of 2007.

23. I have reviewed Exhibit APPLE-1057, “Noninvasive Pulse Oximetry Utilizing Skin Reflectance Photoplethysmography” by Y. Mendelson and B.D. Ochs as published in *IEEE Transactions on Biomedical Engineering* in October of 1988.

24. Attached hereto as Appendix MENDLESON01, is a true and correct digital copy of “Noninvasive Pulse Oximetry Utilizing Skin Reflectance Photoplethysmography” by Y. Mendelson and B.D. Ochs as held by the Carnegie Library of Pittsburgh library. I secured this copy myself from the library’s collection. In comparing Exhibit APPLE-1057 to Appendix MENDLESON01, it is my determination that Exhibit APPLE-1057 is a true and correct copy of “Noninvasive Pulse Oximetry Utilizing Skin Reflectance

Photoplethysmography” by Y. Mendelson and B.D. Ochs as published in *IEEE Transactions on Biomedical Engineering* in October of 1988.

25. Attached hereto as Appendix MENDELSON02 is a true and correct copy of the MARC record for *IEEE Transactions on Biomedical Engineering* as held by the Carnegie Library of Pittsburgh. I secured this record myself from the library’s public catalog. The MARC record contained within Appendix MENDELSON02 accurately describes the title, author, publisher, and ISSN number of *IEEE Transactions on Biomedical Engineering*. The ‘Lib. Has’ field of this record visible on page 2 indicates this collection includes the October 1988 edition of *IEEE Transactions on Biomedical Engineering* containing “Noninvasive Pulse Oximetry Utilizing Skin Reflectance Photoplethysmography”. In comparing Exhibit APPLE-1057 to Appendix MENDELSON02, it is my determination that Exhibit APPLE-1057 is a true and correct copy of “Noninvasive Pulse Oximetry Utilizing Skin Reflectance Photoplethysmography” by Y. Mendelson and B.D. Ochs as published in *IEEE Transactions on Biomedical Engineering* in October of 1988.

26. The 008 field of each MARC record in Appendices MENDELSON02 indicates the date of record creation. The 008 field of Appendix MENDELSON02 indicates the Carnegie Library of Pittsburgh first

cataloged this journal as of November 1, 1975 and held the journal in perpetuity until 2009. Considering this information, it is my determination that *IEEE Transactions on Biomedical Engineering* October 1988 and therefore “Noninvasive Pulse Oximetry Utilizing Skin Reflectance Photoplethysmography” was made available to the public shortly after its initial release in Autumn of 1988.

27. Attached hereto as Appendix IEEE03 is a true and correct copy of a declaration made by Mr. Gordon MacPherson, “Director Board Governance & IP Operations of The Institute of Electrical and Electronics Engineers, Incorporated.” Mr. MacPherson’s declaration Appendix IEEE03 states that the “IEEE publishes and makes available technical articles and standards” as part of its “ordinary course of business,” and that these publications are “made available for public download through the IEEE digital library, IEEE Xplore.”

28. Mr. MacPherson’s declaration includes, as Exhibit A, an article referred to as “Y. Mendelson and B.D. Ochs, ‘Noninvasive pulse oximetry utilizing skin reflectance photoplethysmography’, *IEEE Transactions on Biomedical Engineering*, Vol. 35, Issue 10, October 1988,” which was obtained “through IEEE Xplore, where it is maintained in the ordinary course of IEEE’s



business.” Exhibit A also includes a screen capture of the IEEE Xplore portal page for the above-noted article.

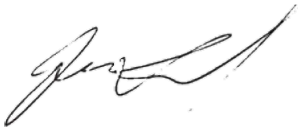
29. In comparing the listed fields in Exhibit A of Appendix IEEE03 to Appendix MENDELSON01 and comparing Exhibit A of Appendix IEEE03 to Exhibit APPLE-1057, it is my determination that Exhibit APPLE-1057 is a true and correct copy of the “Noninvasive Pulse Oximetry Utilizing Skin Reflectance Photoplethysmography” article, and that the copy of the 1988 IEEE publication in the Carnegie Library of Pittsburgh includes the article in Exhibit APPLE-1057.

30. I have been retained on behalf of the Petitioner to provide assistance in the above-illustrated matter in establishing the authenticity and public availability of the documents discussed in this declaration. I am being compensated for my services in this matter at the rate of \$100.00 per hour plus reasonable expenses. My statements are objective, and my compensation does not depend on the outcome of this matter.

31. I declare under penalty of perjury that the foregoing is true and correct. I hereby declare that all statements made herein of my own knowledge are true and that all statements made on information and belief are believed to

be true; and further that these statements were made the knowledge that willful false statements and the like so made are punishable by fine or imprisonment, or both, under Section 1001 of Title 18 of the United States Code.

Dated: 10/25/2021

A handwritten signature in black ink, appearing to read 'Jacob R. Munford', written in a cursive style.

Jacob Robert Munford

**CV**

J. Munford  
Curriculum Vitae

## **Education**

University of Wisconsin-Milwaukee - MS, Library & Information Science, 2009  
Milwaukee, WI

- Coursework included cataloging, metadata, data analysis, library systems, management strategies and collection development.
- Specialized in library advocacy, cataloging and public administration.

Grand Valley State University - BA, English Language & Literature, 2008  
Allendale, MI

- Coursework included linguistics, documentation and literary analysis.
- Minor in political science with a focus in local-level economics and government.

## **Professional Experience**

Researcher / Expert Witness, October 2017 – present  
Freelance ● Pittsburgh, Pennsylvania & Grand Rapids, Michigan

- Material authentication and public accessibility determination. Declarations of authenticity and/or public accessibility provided upon research completion. Experienced with appeals and deposition process.
- Research provided on topics of public library operations, material publication history, digital database services and legacy web resources.
- Past clients include Alston & Bird, Arnold & Porter, Baker Botts, Fish & Richardson, Erise IP, Irell & Manella, O'Melveny & Myers, Perkins-Coie, Pillsbury Winthrop Shaw Pittman and Slayden Grubert Beard.

Library Director, February 2013 - March 2015  
Dowagiac District Library ● Dowagiac, Michigan

- Executive administrator of the Dowagiac District Library. Located in

Southwest Michigan, this library has a service area of 13,000, an annual operating budget of over \$400,000 and total assets of approximately \$1,300,000.

- Developed careful budgeting guidelines to produce a 15% surplus during the 2013-2014 & 2014-2015 fiscal years while being audited.
- Using this budget surplus, oversaw significant library investments including the purchase of property for a future building site, demolition of existing buildings and building renovation projects on the current facility.
- Led the organization and digitization of the library's archival records.
- Served as the public representative for the library, developing business relationships with local school, museum and tribal government entities.
- Developed an objective-based analysis system for measuring library services - including a full collection analysis of the library's 50,000+ circulating items and their records.

November 2010 - January 2013

Librarian & Branch Manager, Anchorage Public Library ● Anchorage, Alaska

- Headed the 2013 Anchorage Reads community reading campaign including event planning, staging public performances and creating marketing materials for mass distribution.
- Co-led the social media department of the library's marketing team, drafting social media guidelines, creating original content and instituting long-term planning via content calendars.
- Developed business relationships with The Boys & Girls Club, Anchorage School District and the US Army to establish summer reading programs for children.

June 2004 - September 2005, September 2006 - October 2013

Library Assistant, Hart Area Public Library

Hart, MI

- Responsible for verifying imported MARC records and original MARC

cataloging for the local-level collection as well as the Michigan Electronic Library.

- Handled OCLC Worldcat interlibrary loan requests & fulfillment via ongoing communication with lending libraries.

## **Professional Involvement**

Alaska Library Association - Anchorage Chapter

- Treasurer, 2012

Library Of Michigan

- Level VII Certification, 2008
- Level II Certification, 2013

Michigan Library Association Annual Conference 2014

- New Directors Conference Panel Member

Southwest Michigan Library Cooperative

- Represented the Dowagiac District Library, 2013-2015

## **Professional Development**

Library Of Michigan Beginning Workshop, May 2008

Petoskey, MI

- Received training in cataloging, local history, collection management, children's literacy and reference service.

Public Library Association Intensive Library Management Training, October 2011

Nashville, TN

- Attended a five-day workshop focused on strategic planning, staff management, statistical analysis, collections and cataloging theory.

Alaska Library Association Annual Conference 2012 - Fairbanks, February 2012

Fairbanks, AK

- Attended seminars on EBSCO advanced search methods, budgeting, cataloging, database usage and marketing.

## **Depositions**

2019 • Fish & Richardson

IPR Petitions of 865 Patent, Apple v. Qualcomm (IPR2018-001281 / 39521-00421IP & IPR2018-01282 / 39521-00421IP2)

2019 • Erise IP

Implicit, LLC v. Netscout Systems, Inc (Civil Action No. 2:18-cv-53-JRG)

2019 • Perkins-Coie

Adobe Inc. v. RAH Color Technologies LLC (Cases IPR2019-00627, IPR2019-00628, IPR2019-00629 and IPR2019-00646)

2020 • O'Melveny & Myers

Maxell, Ltd. v. Apple Inc. (Case 5:19-cv-00036-RWS)

2021 • Pillsbury Winthrop Shaw Pittman LLP

Intel v. SRC (Case IPR2020-1449)

## **Limited Case History & Potential Conflicts**

Alston & Bird

- Nokia (v. Neptune Subsea, Xtera)

Arnold & Porter

- Ivantis (v. Glaukos)

Erise I.P.

- Apple
  - v. Future Link Systems (IPRs 6317804, 6622108, 6807505, and 7917680)
  - v. INVT
  - v. Navblazer LLC (Case No. IPR2020-01253)

v. Qualcomm (IPR2018-001281, 39521-00421IP, IPR2018-01282, 39521-00421IP2)

v. Quest Nettech Corp, Wynn Technologies (Case No. IPR2019-00XXX, RE. Patent Re38137)

- Fanduel (v CGT)
  
- Garmin (v. Phillips North America LLC, Case No. 2:19-cv-6301-AB-KS Central District of California)
  
- Netscout
  - v. Longhorn HD LLC)
  - v. Implicit, LLC (Civil Action No. 2:18-cv-53-JRG)
  
- Sony Interactive Entertainment LLC
  - v. Bot M8 LLC
  - v. Infernal Technology LLC
  
- Unified Patents (v GE Video Compression, Civil Action No. 2:19-cv-248)

#### Fish & Richardson

- Apple
  - v. LBS Innovations
  
  - v. Masimo (IPR 50095-0012IP1, 50095-0012IP2, 50095-0013IP1, 50095-0013IP2, 50095-0006IP1)
  
  - v. Neonode
  
  - v. Qualcomm (IPR2018-001281, 39521-00421IP, IPR2018-01282, 39521-00421IP2)
  
- Dish Network
  - v. Realtime Adaptive Streaming, Case No 1:17-CV-02097-RBJ)



v. TQ Delta LLC

- Huawei (IPR 76933211)
- Kianxis
- LG Electronics (v. Bell Northern Research LLC, Case No. 3:18-cv-2864-CAB-BLM)
- Metaswitch
- MLC Intellectual Property (v. MicronTech, Case No. 3:14-cv-03657-SI)
- Realtek Semiconductor
- Quectel
- Samsung (v. Bell Northern Research, Civil Action No. 2:19-cv-00286-JRG)
- Texas Instruments

Irell & Manella

- Curium

O'Melveny & Myers

- Apple (v. Maxell, Case 5:19-cv-00036-RWS)

Perkins-Coie

- TCL Industries (v. Koninklijke Philips NV, PTAB Case Nos. IPR2021-00495, IPR2021-00496, and IPR2021-00497)

Pillsbury Winthrop Shaw Pittman

- Intel (v. FG SRC LLC, Case No. 6:20-cv-00315 W.D. Tex)

**DING01**

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# Physics in Medicine & Biology

Volume 51 Number 6 21 March 2006

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# Physics in Medicine and Biology

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# Physics in Medicine and Biology

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## Refractive indices of human skin tissues at eight wavelengths and estimated dispersion relations between 300 and 1600 nm

Huafeng Ding<sup>1</sup>, Jun Q Lu<sup>1</sup>, William A Wooden<sup>2</sup>, Peter J Kragel<sup>3</sup> and Xin-Hua Hu<sup>1</sup>

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### Abstract

The refractive index of human skin tissues is an important parameter in characterizing the optical response of the skin. We extended a previously developed method of coherent reflectance curve measurement to determine the *in vitro* values of the complex refractive indices of epidermal and dermal tissues from fresh human skin samples at eight wavelengths between 325 and 1557 nm. Based on these results, dispersion relations of the real refractive index have been obtained and compared in the same spectral region.

(Some figures in this article are in colour only in the electronic version)

### 1. Introduction

Understanding the response of the skin to optical radiation is essential to the dermatological applications of photomedicine. Among various skin optical parameters, refractive index is an important one. At the microscopic scales ranging from 1 to 10  $\mu\text{m}$ , refractive index variation causes light scattering which can be understood by direct solution of the Maxwell equations within the framework of classical electrodynamics for simple shaped particles (Bohren and Huffman 1983, Ma *et al* 2003a) and for biological cells (Lu *et al* 2005). For highly turbid tissues of human skin with sizes of 1 mm or larger, modelling of tissue optics based on the electrodynamic theory is very difficult, and the real refractive index and scattering parameters are often treated as independent parameters within the frameworks of effective medium theory and radiative transfer theory, respectively. For example, in the widely used method of Monte Carlo simulation of light distribution in biological tissues, photon interaction with an interface



between tissue regions of different refractive indices is described according to the Fresnel equations which require the index as the key parameter (van Gemert *et al* 1989, Wang *et al* 1995, Lu *et al* 2000). Furthermore, accurate modelling of measured light signals near a skin surface requires the refractive indices of skin tissues to account for the redistribution of light due to the index mismatch at the surface (Lu *et al* 2000, Meglinsky and Matcher 2001, Bartlett and Jiang 2001, Ma *et al* 2003b, 2005).

Determination of the refractive indices of the human skin tissues, however, presents challenges because of their highly turbid natures. In transparent or absorbing media such as the aqueous solutions with molecular solutes, propagation of light is dominated by its coherent component. The reflection and refraction of a light beam, as it passes through an interface between two media of different refractive indices, are described with the Fresnel equations on the basis of an effective medium theory. In contrast, light propagation in a turbid medium produces a scattering component that becomes increasingly dominant as the light penetrates into the medium. This feature of interaction often precludes the use of refraction method to determine the refractive index of biological tissue samples where uniform and thin samples are very difficult to obtain. Recently, we developed an automated reflectometer system for determining the refractive index of a turbid sample by measurement of its coherent reflectance  $R$  versus the incident angle  $\theta$  without the need to section skin tissues (Ding *et al* 2005). Here we report complex refractive indices of fresh human skin tissues determined by nonlinear regression of  $R(\theta)$  with the Fresnel equations. The complex refractive index has been obtained at eight wavelengths between 325 and 1557 nm for both the epidermis and dermis tissues. With these data we investigated various dispersion schemes for interpolation of the index data at other wavelengths in this spectral region.

## 2. Methods

Fresh skin tissue patches were obtained from the patients undergoing abdominoplasty procedures at the plastic surgery clinic of the Brody School of Medicine, East Carolina University (ECU). A study protocol approved by the Institutional Review Board of ECU was strictly followed and a consent form was signed by each participating patient before the surgery. We obtained one skin tissue patch from each of the 12 female patients with ages between 27 and 63 years old; 10 are Caucasian and 2 are African Americans, with the skin data compiled in table 1. Each skin patch was stored in a bucket on crushed ice ( $\sim 2^\circ\text{C}$ ) inside a refrigerator immediately after surgery. Samples with sizes of about  $1\text{ cm} \times 1\text{ cm}$  were prepared by removing the hair on the skin surface with scissors and subcutaneous fat tissue with a razor blade and warming the skin to a room temperature of about  $22^\circ\text{C}$  with 0.9% saline drops. Care was taken to preserve the stratum corneum layer of the skin epidermis. The skin sample was pressed against the base of the prism with a pistol pressurized by a nitrogen gas cylinder to maintain good contact between the sample and the prism. The periphery of the tissue sample between the pistol and prism base was sealed with plastic tape to prevent sample dehydration during the measurement. By pressing either the epidermis or dermis side of the skin sample against the prism base, the coherent reflectance curves of skin epidermis or dermis were measured, respectively. All reflectance curve measurements were performed at the room temperature within 30 h after the abdominoplasty procedure.

An automated reflectometer system has been designed and constructed to measure the coherent reflectance as a function of incident angle. Compared to other approaches of index determination based on fibre insertion and OCT (Bolin *et al* 1989, Tearney *et al* 1995, Knuttel and Boehlau-Godau 2000), this method has the combined benefits of high accuracy, wide spectral capability and instrumentation simplicity. The system has been described in detail

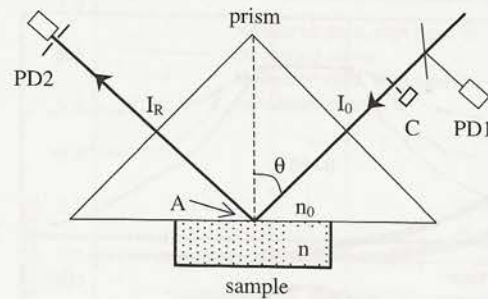


Figure 1. The schematic of the reflectometer system.

Table 1. The human skin sample data.

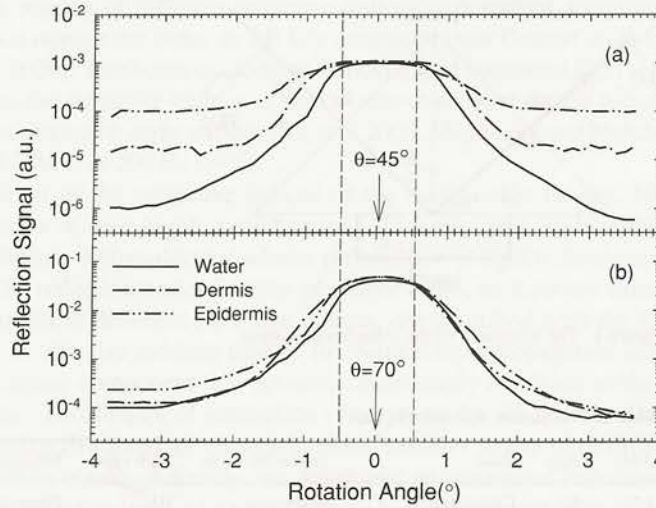
ID no.	Age	Race	Tissue location	Skin type	Measurement
1	42	Caucasian	Abdomen	III	Pressure dependence
2	40	Caucasian	Abdomen	I	633 nm, 532 nm
3	27	African American	Abdomen	V	442 nm
4	63	Caucasian	Abdomen	II	1064, 850 nm
5	56	Caucasian	Abdomen	II	325, 1550 nm
6	54	Caucasian	Arm	II	1310, 633 nm
7	34	Caucasian	Abdomen	II	1064, 325 nm
8 <sup>a</sup>	55	Caucasian	Abdomen	I	532, 633 nm
9 <sup>a</sup>	49	Caucasian	Abdomen	III	442, 1310 nm
10 <sup>a</sup>	41	Caucasian	Abdomen	II	850, 1550 nm
11 <sup>a</sup>	39	African American	Abdomen	V	532, diffuse reflection
12	44	Caucasian	Abdomen	III	Pressure dependence

<sup>a</sup> The skin structures of the samples from these patients have been examined through histology.

elsewhere (Ding *et al* 2005). Briefly, a right-angle glass prism was used to interface with a skin sample on the prism base and a linear polarized laser beam was propagated through one side surface as the incident beam on the prism–sample interface at an incident angle of  $\theta$ . The coherent reflectance  $R$  of the laser beam was measured at the angle of specular reflection by a photodiode of either Si or GaAs, depending on the light wavelength. Two coherent reflectance curves,  $R_s(\theta)$  or  $R_p(\theta)$ , have actually been measured for each sample with either s- or p-polarized incident beam, respectively. The incident angle  $\theta$  can be varied between  $48^\circ$  and  $80^\circ$  with a stepsize of  $0.125^\circ$  and resolution of  $0.006^\circ$ . A schematic diagram of the optical setup is presented in figure 1. The powers of the incident and reflected beams were measured by two identical photodiodes and the effect of the reflection loss at the side surfaces of the prism was removed to determine the coherent reflectance ( $R_s$  or  $R_p$ ). To reduce the contribution of diffuse reflection to reflection signal, a pinhole of 2 mm diameter was used in front of the photodiode, resulting in an angular range of  $1.74 \times 10^{-2}$  rad or about  $1.00^\circ$  in the measurement. The incident laser beam, modulated at 370 Hz for lock-in detection, was produced by one of seven continuous-wave (cw) lasers at one of eight wavelengths:  $\lambda = 325, 442, 532, 633, 850, 1064, 1310$  and  $1557$  nm. The diameter of the beam was set to between 1 and 2 mm with the incident beam power adjusted to be about  $1 \mu\text{W}$ .

The measured coherent reflectance curves have been fitted by the calculated values,  $\tilde{R}_s(\theta)$  and  $\tilde{R}_p(\theta)$ , according to the Fresnel equations. The fitting requires the assumed value of the





**Figure 2.** The reflection signal versus rotation angle of the detector at the incident angle of (a)  $\theta = 45^\circ$ ; (b)  $\theta = 70^\circ$  with a s-polarized beam at  $\lambda = 633$  nm for deionized water, the epidermis and dermis of one skin sample. The error bars of about  $\pm 5\%$  were removed for clear view and the two dashed lines indicate the angular acceptance range of the aperture in front of the photodiode.

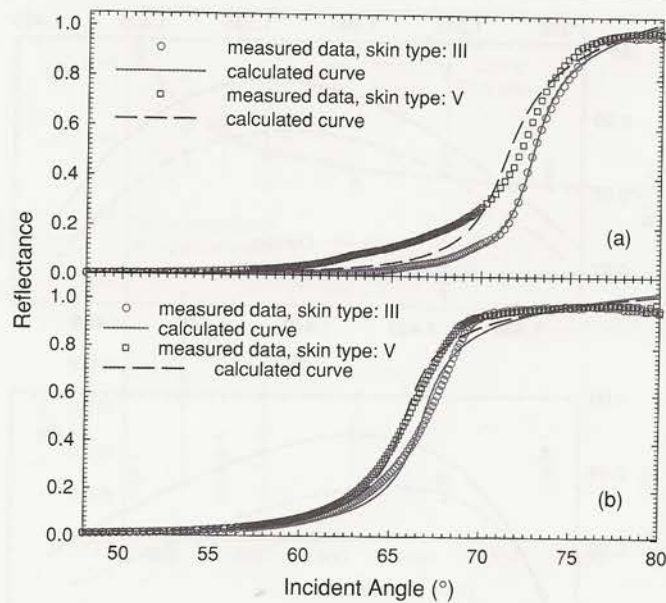
complex refractive index of the turbid sample,  $n = n_r + in_i$ , and the known refractive index,  $n_0$ , of the prism. Therefore, the index  $n$  was inversely determined using an iteration process to achieve least-squared difference between the calculated and measured curves (Ding *et al* 2005). The consistency between the measured and calculated coherent reflectance curves is described by a coefficient of determination,  $R^2$ , defined as

$$R^2 = 1 - \frac{\sum_{i=1}^N (R_i - \tilde{R}_i)^2}{\sum_{i=1}^N (R_i - \bar{R})^2}, \quad (1)$$

where  $R_i$  and  $\tilde{R}_i$  denote the measured and calculated reflectances at the  $i$ th angle of incidence  $\theta_i$ , respectively, and  $\bar{R}$  is the mean value of measured reflectance over  $N$  values of  $\theta$ . The  $R^2$  value ranges between 0 and 1 with  $R^2 = 1$  for a perfect fit. The system was calibrated before measurements of each sample by comparing the measured real refractive index of deionized water with the published value at the wavelength of measurements (Hale and Querry 1973). From the water data, the experimental error in determination of the real refractive index  $n_r$  of transparent samples by the reflectometer system was found to be about  $\delta n_r = \pm 0.002$ .

### 3. Results

To ensure that the reflection signal is dominated by its coherent component, we measured the angular distribution of the reflected beam around a specular reflection angle at two positions ( $\theta = 45^\circ$  or  $75^\circ$ ) with  $\lambda = 633$  nm. Similar data with deionized water were used as the baseline and all are plotted in figure 2. These results demonstrate that the contribution of the diffusely reflected light to the coherent reflectance signal is negligible within the  $1^\circ$  angular range defined by the photodiode aperture (indicated by the dashed lines in figure 2). Two typical sets of coherent reflectance curves from the epidermis and dermis sides of the skin



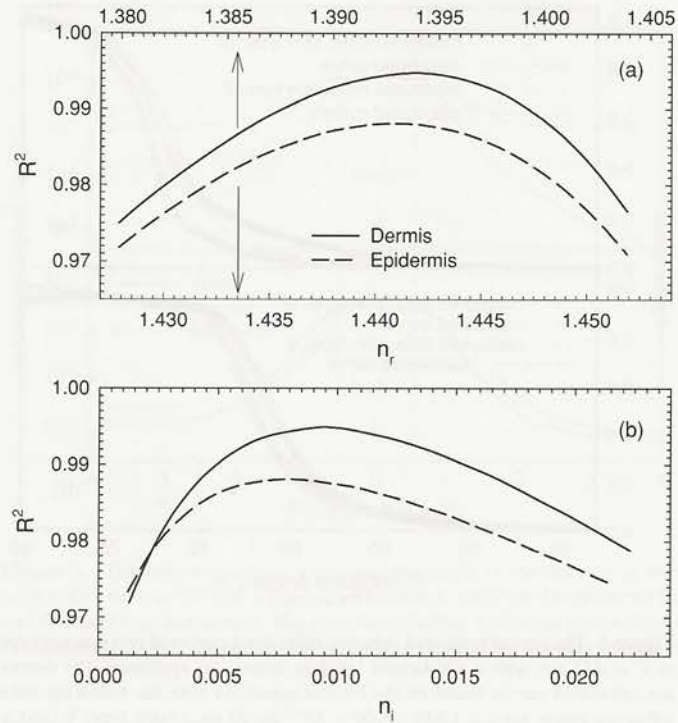
**Figure 3.** The typical measured coherent reflectance curves of two skin samples from two patients at  $\lambda = 442$  nm with a s-polarized incident beam: (a) epidermis, (b) dermis. The solid lines are calculated curves based on the Fresnel equations with the following values of the complex refractive index: (a)  $n = 1.445 + i1.00 \times 10^{-2}$  for ID no. 3 (skin type: V) and  $n = 1.458 + i8.34 \times 10^{-3}$  for ID no. 9 (skin type: III); (b)  $n = 1.394 + i9.30 \times 10^{-3}$  for ID no. 3 and  $n = 1.404 + i9.20 \times 10^{-3}$  for ID no. 9.

samples of two patients with different skin types are presented in figure 3 together with the fitted curves based on the Fresnel equations. From these results one can see that the agreement between the measured and calculated reflectance curves varies from sample to sample and is gauged by the coefficient of determination  $R^2$ . To determine the sensitivity of the refractive index on the nonlinear regression, we analysed the relation between  $R^2$  and  $n$  with selected data of the coherent reflection curves and typical results are presented in figure 4. On the basis of the standard deviations in the distribution of  $R^2$  values, we estimated that the uncertainty in obtaining the real and imaginary is about  $\Delta n_r = \pm 0.006$  and  $\Delta n_i = \pm 0.005$ , respectively, for the turbid tissues of both the epidermis and dermis.

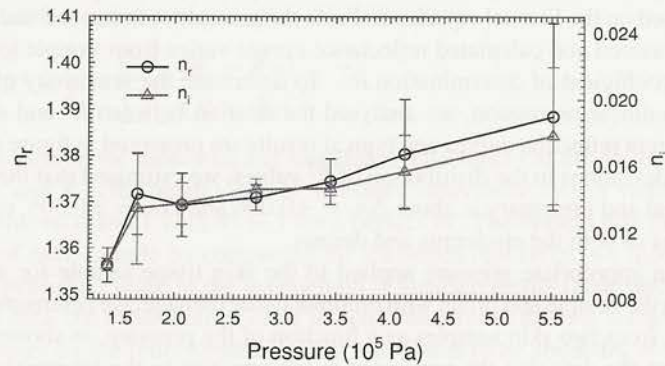
To select an appropriate pressure applied to the skin tissue sample for achieving good contact between the sample and prism with minimal tissue damage, we determined the complex refractive index from two skin samples as a function of the pressure, as shown in figure 5. It can be seen from the data that the real index is not sensitive to the air pressure set between  $2 \times 10^5$  and  $5 \times 10^5$  Pa. Signs of damage to the skin samples became visible when the pressure was increased to above  $4 \times 10^5$  Pa, which included the fast dehydration of the tissue samples and significant reduction in the dermis layer thickness. On the basis of these results, all subsequent measurements of coherent reflectance of skin tissue samples were carried out at a fixed pressure of  $2.06 \times 10^5$  Pa (30 psi or 2.0 atm) to minimize possible structural change in the skin tissues.

At each wavelength, 8 or 12 skin samples were used to measure the coherent reflectance curves of  $R_s(\theta)$  and  $R_p(\theta)$  with 4 from one patient. Half of the samples were measured with the epidermis side in contact with the prism base and half with the dermis side. The measurement of  $R_s(\theta)$  and  $R_p(\theta)$  was repeated three times on the same skin sample and thus the data set



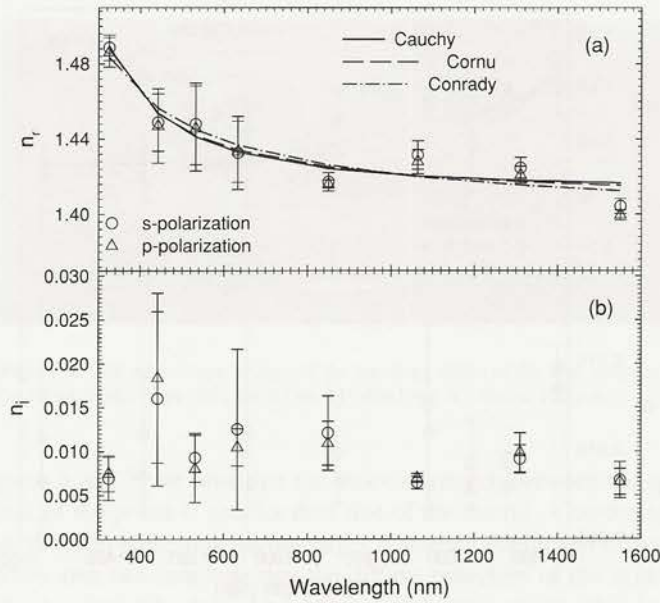


**Figure 4.** The dependence of the coefficient of determination  $R^2$  on different choices of  $n_i$  for a coherent reflectance curve measured from the epidermis and dermis sides of a skin sample at  $\lambda = 442$  nm.



**Figure 5.** The average real and imaginary refractive index versus sample pressure determined from the dermis of two skin samples.

at each wavelength consists of 12 or 18 curves with an incident beam of s- or p-polarization. Nonlinear regression to the coherent reflectance curve data by the Fresnel equations was done individually to obtain the complex refractive index from each measurement. The coefficient of determination  $R^2$  ranges from 0.960 to 0.999 for the data from the measurement of the epidermis side and from 0.978 to 0.998 for the dermis side. The mean values and standard deviations of the complex refractive index have been calculated at each wavelength from the



**Figure 6.** The (a) real and (b) imaginary refractive indices of human skin epidermis versus wavelength. Each data point and associated error bar is the mean and standard deviation obtained from 12 or 18 measurements of 4 or 6 skin samples. The lines in (a) are based on the dispersion equations.

data sets. These results are plotted as a function of wavelength in figure 6 for the epidermis and figure 7 for the dermis.

We investigated various dispersion schemes to identify appropriate ones for calculation of real refractive index of human skin tissues at wavelengths between 300 and 1600 nm based on our index data at eight wavelengths. Among those reported on the index data of ocular tissues (Kroger 1992, Atchison and Smith 2005), we selected three schemes to fit to our data: the Cauchy dispersion equation

$$n_r = A + \frac{B}{\lambda^2} + \frac{C}{\lambda^4}, \quad (2)$$

the Cornu equation

$$n_r = A + \frac{B}{(\lambda - C)} \quad (3)$$

and the Conrady equation

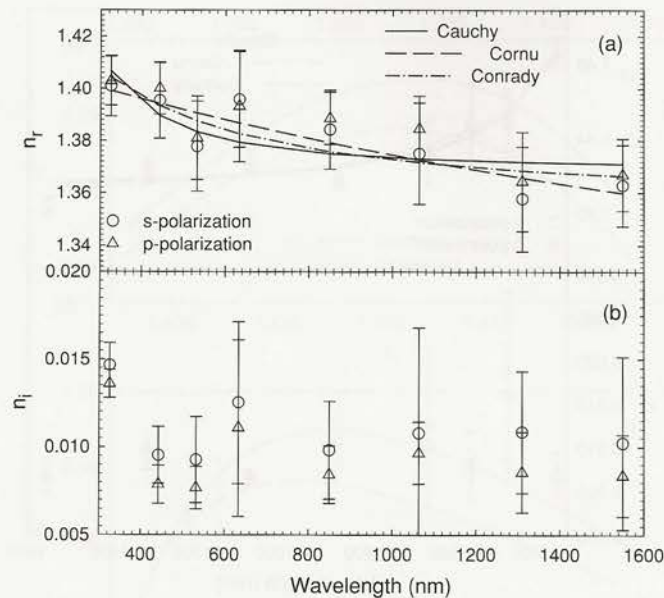
$$n_r = A + \frac{B}{\lambda} + \frac{C}{\lambda^{3.5}}. \quad (4)$$

The coefficients of each dispersion scheme determined with the least-squares principle from our index data are given in table 2.

#### 4. Discussion

Refractive index plays an important role in characterization of the biological tissues' response to optical illumination, particularly for tissues of heterogeneous composition such as the layered skin tissues (Tuchin 2005). The real refractive index not only influences optical





**Figure 7.** The (a) real and (b) imaginary refractive indices of human skin dermis versus wavelength. Each data point and associated error bar is the mean and standard deviation obtained from 12 or 18 measurements of 4 or 6 skin samples. The lines in (a) are based on the dispersion equations.

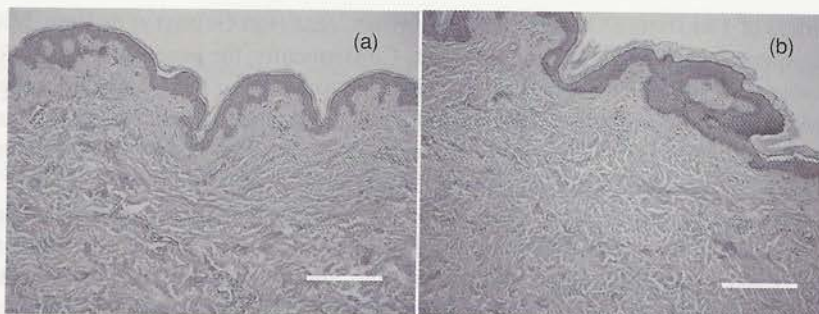
**Table 2.** The coefficients of different dispersion equations<sup>a</sup>.

Dispersion equation	A	B	C
Cauchy	1.3696	$3.9168 \times 10^3$	$2.5588 \times 10^3$
Cornu	1.2573	$4.5383 \times 10^2$	$2.8745 \times 10^3$
Conrady	1.3549	$1.7899 \times 10$	$-3.5938 \times 10^6$

<sup>a</sup> These coefficients were obtained on the basis of equations (2)–(4) with wavelength in the unit of nanometres.

pathlengths of light propagating in tissues that can be determined by the methods of phase or time-resolved spectroscopy (Duncan *et al* 1995, Zhao *et al* 2002) but also affects the light measurement outside the tissues due to the index mismatch at the boundaries (Meglinsky and Matcher 2001, Bartlett and Jiang 2001, Ma *et al* 2003b, 2005). With an automated reflectometer system, we have extended the technique of measuring the coherent reflectance curve (Meeten and North 1995, Ding *et al* 2005) to obtain the complex refractive indices of fresh human skin tissues. These results, therefore, are of interest to researchers wishing to conduct quantitative optical studies involving index-mismatched interfaces in the human skin.

The measurement of coherent reflectance was validated by confirming the dominance of the coherent reflection over the diffused one at the specular reflection angle for both of the epidermis and dermis sides of the skin samples, as shown in figure 2. Diffuse reflection occurs mainly outside the angular range defined by the aperture of the photodiode PD2 (see figure 1) and decreases towards the baseline data of water for large incident angles as  $\theta$  approaches  $80^\circ$ . The diffusely reflected light originates from two sources: the rough tissue surface mismatched optically with the prism glass and the tissue bulk. From the index data



**Figure 8.** The microscope images of the histology slides of the skin samples from two patients: (a) ID no. 9 (skin type: III); (b) ID no. 11 (skin type: V). Bar = 100  $\mu\text{m}$ .

presented in figures 6 and 7, we note that the index mismatch between the epidermis tissue and the BK7 glass of the prism is smaller than that of the dermis. Combining this fact with the knowledge of the skin epidermis having larger scattering coefficient than the dermis (van Gemert *et al* 1989), one can conclude that the diffuse reflection of the skin tissues seen in figure 2 should be dominated by the light scattering in the tissue bulk. This is consistent with our previous results on comparison of the diffuse reflection between the porcine skin tissues with rough surfaces and the intralipid solution samples with smooth surface (Ding *et al* 2005).

The human skin has a layered structure with two primary layers of epidermis and dermis, both are beneath the superficial layer of the epidermis or the stratum corneum (sc). We examined the tissue structures by preparing histological slides of the skin tissue samples from 4 patients (with ID no. from 8 to 11, see table 1) with standard H&E staining. Two examples of skin slides are shown in figure 8 with one from a Caucasian and another from an African American patient. It can be seen that the sc layer is less than 10  $\mu\text{m}$  in thickness, as expected, with the thickness of epidermis ranging from about 30 to 80  $\mu\text{m}$ . We further verified that the sc layer has no significant effect on the refractive index determination by comparing the index values from samples with and without the sc layer prepared from fresh porcine skin tissues. The sc layer was removed from the epidermal side of the porcine tissue samples using a tape-stripping method (Beisson *et al* 2001) without heating the samples. The real refractive index  $n_r$  of the epidermis at the wavelengths of 442 nm and 1064 nm was found to be the same within the experimental errors between the samples with and without the sc layer. These results demonstrated that the sc layer has no significant effect on the real refractive index of the skin epidermis because of its small thickness in comparison with the penetration depth (Everett *et al* 1966, van Gemert *et al* 1989), see also the discussion below. Another point worth noting is the effect of the melanin on the refractive index of the skin tissues. As can be seen from figure 8, many basal keratinocytes containing melanin pigment are visible near the epidermal–dermal junction in patient no. 11 (skin type: V), while little pigment exists in patient no. 9 (skin type III). The high density of melanin in the type V skins appears only to affect significantly the imaginary refractive index of epidermis, as shown in figure 3.

A general model of refractive index for a dense and turbid medium remains an open question. But according to the existing models of effective medium for absorbing or dilute turbid media (Ballenegger and Weber 1999, Barrera and Garcia-Valenzuela 2003) the refractive index determined from a coherent reflectance curve should relate to the medium's optical response from over at least the full depth of penetration of the coherent component of the incident wave in the medium. The total attenuation coefficients, as the sum of the scattering and absorption coefficients, for both the skin epidermis and dermis are expected to be on the



orders of 1 to 10 mm<sup>-1</sup> based on the published data (van Gemert *et al* 1989, Ma *et al* 2005) in the spectral region from 300 to 1600 nm. Consequently, the penetration depth for the coherent component should be about a few hundred micrometres or less. Therefore, one would expect the tissue response of the first 100  $\mu$ m layer to dominate the coherent reflectance and thus the value of the real refractive index. This conclusion is supported by the wavelength correlation of the real refractive index determined from the epidermis and dermis sides of the skin tissue samples. The correlation coefficient of wavelength dependence of the real refractive index was found to be  $r_{\text{corr}} = 0.99$  between the index determined with s- and p-polarized beam for the epidermis and  $r_{\text{corr}} = 0.95$  for the dermis. The values of  $r_{\text{corr}}$  decrease drastically to 0.057 and 0.065 between the index of epidermis and dermis measured with the s- and p-polarized beam, respectively.

Within the previously discussed uncertainty on the real refractive index, our *in vitro* results of  $n_r$  at about  $\lambda = 1310$  nm agree with those determined *in vivo* from human skin epidermis (averaged over the sc and other sub-layers) by the OCT method (Tearney *et al* 1995, Knuttel and Boehlau-Godau 2000). The  $n_r$  of dermis, however, is smaller than the above reported *in vivo* results: 1.36 versus 1.41 at about  $\lambda = 1310$  nm. The wavelength dependence is similar to the  $n_r$  of bovine and porcine muscle tissues in the visible region (Bolin *et al* 1989, Li and Xie 1996). To extend the use of our real refractive index data on a limited number of wavelengths, we have tested different dispersion schemes based on the equations by Cauchy, Cornu and Conrady. From figures 6 and 7, it is clear that these relations are close to each other and all fit to data fairly well for the dermis and very well for the epidermis. Therefore, these equations may be used to estimate the values of the real refractive indices of human skin tissues with the coefficients given in table 2 between 300 and 1600 nm. These estimations should be further improved as the refractive index becomes available at an increased number of wavelengths in this spectral region.

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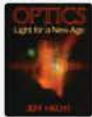
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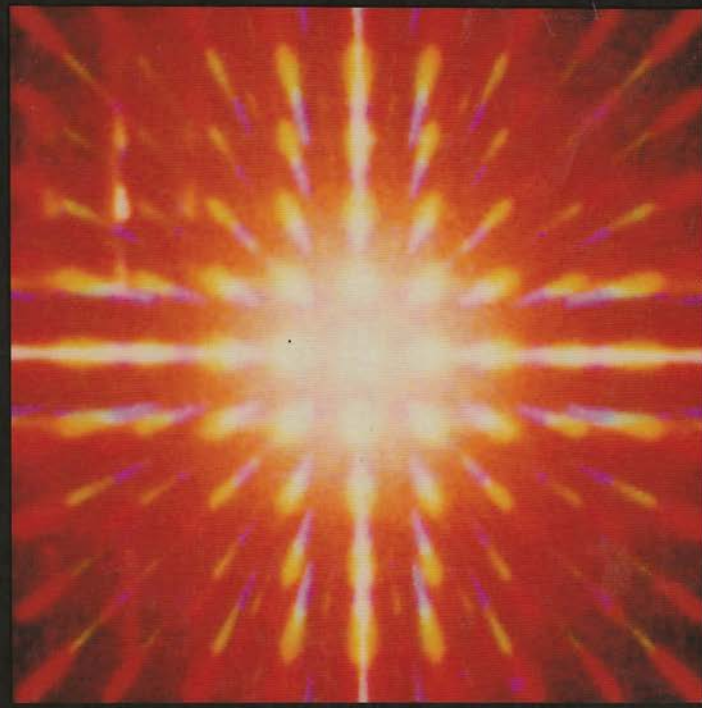
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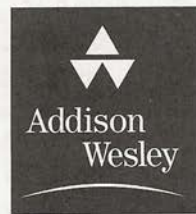


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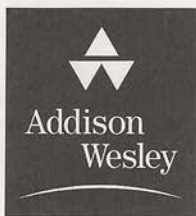
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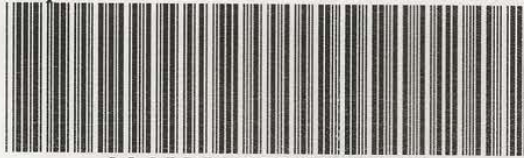
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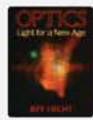
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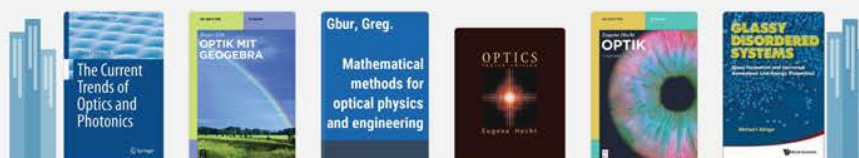
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## 1. Introduction

At present plastic materials find wide application in consumer and high quality optics. Optical plastics (OPs) are used mainly in the visible (VIS) and near-infrared (NIR) spectral regions from 400 nm to 1100 nm. Success in application of OPs depends on knowledge of their optical refraction, transmission, birefringence, haze and homogeneity [1]. The optical properties of polymers are in details considered in [2]. Chromatic dispersion is an important characteristic in the design of optical systems and devices. However, measurements of refractive indices are usually realized at several selected wavelengths. Determination of more extensive refractometric data is possible using dispersion formulae [3,4].

The measuring methods for determination of OPs' refractometric characteristic are quite different. The refractive indices of transparent polymers can be obtained using the Federal Test Method Standard [5] in which the Abbe refractometer is applied. It operates with a white light source and Amici prisms as colour compensators. The refractive index value for the sodium *D*-line can be read directly from the instrument. However, the measuring accuracy is not acceptable for modern optical design projects. Furthermore, determination of refractive indices values can not be realized at different wavelengths. Utilization of Zeiss Pulfrich refractometer (PR2) is possible too [6,7]. We have measured the refractive indices in the VIS light using the PR2 instrument with its V-type prism [8] and additional goniometric set-up was applied for the entire

VIS and NIR regions. The obtained refractive values were compared with the data from Glass catalogues of OSLO [9], ZEMAX [10] and Code V [11]. Laser refractometric measurements of a number of OP specimens have been also accomplished using a He–Ne laser source with 632.8 nm emission wavelength [12].

Most widely used OPs are thermoplastics as polymethyl methacrylate (PMMA), polystyrene (PS), polycarbonate (PC), methyl methacrylate styrene copolymer (NAS), styrene acrylonitrile (SAN) and methylpentene (TPX) [1,13]. The only thermosetting plastic for optical applications is allyl diglycol carbonate (CR-39) [1]. The available catalogue data for refractive indices and dispersion characteristics of OPs is yet scanty. Useful data of commonly used transparent polymers is presented in Refs. [1,13]. The refractive indices of the principal OPs are included in many patents [14–17] and available online web-pages [18,19]. Companies producing trade-marks of optical polymers provide information on their refractometric and dispersive characteristics [20]. In comparison with glass, OPs have a restricted range of refractive indices and dispersion. The magnitude of the refractive index  $n_D$  at the sodium *D*-line (589.3 nm) usually varies from 1.47 to 1.59 [1,21]. The Abbe number of OPs is in the range from about 100 to a little less than 20 [1,21]. However, there are some differences in the reported refractometric data. For example, the refractive indices values of such a popular material as PMMA are different in references [15,16].

Recently it has been noticed considerable interest in the development of OP materials. Plastics replace glasses in products as objectives, lens arrays, aspheric and ophthalmic lenses, displays and lighting fixtures, windows, internally illuminated outdoor signs and skylights [21,22]. The improvement of manufacturing processes makes possible the utilization of OPs in medicine, military optics, sensors

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and communications [21–24]. Chemical companies produce various trade-marks of OPs as NAS-21 Novacor<sup>®</sup>, CTE-Richardson<sup>®</sup>, Zeonex<sup>®</sup>, Optorez<sup>®</sup>, Bayer<sup>®</sup>, etc. but some of them have close dispersion data. Our previous refractometric measurements [7] show for example that Optorez 1330<sup>®</sup> and S – low Styrene<sup>®</sup> have similar dispersive properties, Zeonex E48R<sup>®</sup> and COC<sup>®</sup> are equal, and Bayer<sup>®</sup> is a PC-type plastic material. Using some new optical materials the designers can improve the performance and balance the production expenses [13]. It seems that OPs can be both a low-cost alternative to glass and an option that provides more degrees of freedom for product and optical design.

In this paper we consider the refractometric and dispersion properties of OP materials in the region of normal dispersion. We have examined various types of OPs including the principal, selected trade marks and some control samples of polymers.

## 2. Theoretical analysis of dispersion

The interaction between the electromagnetic wave and the medium (the refractive index  $n$ , respectively) depends on its density and individual properties of the molecules on one hand, and the radiation wavelength on the other hand. A characteristic of the materials is the ratio, named a specific refraction  $r = f(n)/\rho$ , where  $\rho$  is the density of the substance and  $f(n)$  is a function of the refractive index. The product of the specific refraction and molar mass is the molar refraction  $R$  often used in practice. Considering the effective electric field acting on a molecule in a polarizable medium, Lorentz and Lorenz have formulated theoretically [2] the molar refraction  $R_{LL}$  as

$$R_{LL} = \frac{n^2 - 1}{n^2 + 2} \frac{M}{\rho} = \frac{1}{3} N_A \alpha, \quad (2.1)$$

where  $M$  is the molar mass,  $\alpha$  is the molecular polarisability and  $N_A$  is the Avogadro's number. The molar refraction is a material characteristic of the refraction properties which is independent on the density, temperature and physical state. This equation is valid for isotropic materials in cases of ultra-violet and visible illumination, where the electron polarisability is essential. For organic liquids Gladstone and Dale have obtained that at standard wavelengths the ratio  $(n - 1)/\rho$  is a material characteristic constant and the molar refraction is:

$$R_{GD} = \frac{M}{\rho} (n - 1). \quad (2.2)$$

Another correlation between the chemical structure of electrically insulating materials and the refractive index has been proposed by Vogel [2]:

$$R_V = nM. \quad (2.3)$$

It is known that the refraction is an additive value. Therefore, the refractive index can be estimated using Eqs. (2.1), (2.2) and (2.3), if the structure of the compounds is known. Goedhart estimated the molar refraction values

of 43 functional groups, using an extensive regression analysis [2]. With his group contributions the refractive index can be predicted with an average standard deviation of about 0.4%.

Variation of the refractive index with respect to the wavelength depends also on the structure of the substance. In the spectral region of transmission the refractive index of the materials reduces towards longer wavelengths. Close to the absorption bands the refractive index enhances with the increase of the wavelength. The major OPs absorb in the blue portion of the visible spectrum and have some energy absorption at wavelengths of 900 nm, 1150 nm and 1350 nm in the NIR region [21]. At longer wavelengths high transmission is possible only in very thin sections of the material (0.022 mm) [1]. OPs become totally opaque at about 2100 nm [21]. The transmission characteristics depend strongly on quenching the material from temperature of about 213 °C to 147 °C in 60 s or less to prevent crystallization [1]. At wavelengths within the absorption bands intramolecular oscillations appear and the bond lengths and valence angles of molecules are altered. The more complex chemical structures of polymers increase the number of absorption bands and therefore influence the dependence of the refractive index on the wavelength.

There are several formulae in literature approximating the dispersion of optical materials. Most popular among them are the Hezberger's experimental formula [3], Cauchy's and Sellmeier's equations [4]. Only the last one, however, has physical ground. Dispersion can be explained by applying the classical electromagnetic theory to the molecular structure of the medium. Sellmeier has considered the substance as a system of elastically bounded particles with natural angular frequency  $\omega_{0i}$  ( $i = 1 \dots k$  – consecutive number of a single oscillator). The amplitude of the oscillations of bound charges forced by the electromagnetic wave increases at resonance frequency. According to Sellmeier's theory the well-known dependence of the refractive index on the wavelength is obtained:

$$n^2(\lambda) = 1 + \sum_{i=0}^k \frac{A_i \lambda^2}{\lambda^2 - \lambda_{0i}^2}, \quad (2.4)$$

where  $A_i$  are constants proportional to the number of oscillators with natural wavelengths  $\lambda_{0i}$  per unit volume.

An important characteristic of the dispersion properties of optical materials is their Abbe numbers, which usually decrease as the refractive indices increase. The USA standard Abbe number  $v_d$  is defined by the following ratio:

$$v_d = \frac{n_d - 1}{n_F - n_C}. \quad (2.5)$$

The difference  $(n_F - n_C)$ , named a principal dispersion, involves refractive indices  $n_F$  and  $n_C$  at the blue hydrogen  $F$ -line (486.13 nm) and red hydrogen  $C$ -line (656.27 nm). The Abbe number  $v_{804}$  which is a measure of partial dispersion in the NIR range is given by the equation

$$v_{804} = \frac{n_{804} - 1}{n_{703} - n_{1052}}, \quad (2.6)$$

where  $n_{703}$ ,  $n_{804}$  and  $n_{1052}$  are the refractive indices at wavelengths 703 nm, 804 nm and 1052 nm, respectively. In Europe and Japan the Abbe number is defined according to the green mercury  $e$ -line (546.07 nm) as:

$$v_e = \frac{n_e - 1}{n_{F'} - n_{C'}}, \quad (2.7)$$

where  $n_{F'}$  and  $n_{C'}$  – refractive indices at the cadmium blue  $F'$ -line and red  $C'$ -line. Materials with low refractive indices usually have low dispersion behaviour and therefore a high Abbe number.

### 3. Measurement of the indices of refraction

In this study we apply the measuring method described in details in our paper [7]. The OPs' indices of refraction were measured with the aid of the Carl Zeiss Jena Pulfrich-Refractometer PR2 [8] in the visible spectral region at six standard spectral lines: green  $e$ -line 546.07 nm, blue  $g$ -line 435.83 nm, yellow  $d$ -line 587.56 nm, red  $r$ -line 706.52 nm, blue  $F$ -line 486.13 nm and red  $C$ -line 656.27 nm. We have chosen the  $V$ -type SF3 glass prism (VoF3 prism) which usually used for measuring liquids since the standard total internal reflection prism requires thick cubic of OP samples with satisfactory polished surfaces to observe a contrast image of the dividing border between the light and dark field in the eyepiece of the instrument. Furthermore, the standard prism does not avoid evaporating of the water contacting solutions and the precision of monitoring decreases at the end of the measuring series.

The examined OP specimens were prepared as injection moulded plates with thickness varying from 2.54 mm to 5.1 mm, except for the control samples which were cubic. Two mutually perpendicular surfaces of the samples were well polished to obtain a good refractometric data. Measuring temperature of 20 °C is maintained by a thermostat and temperature regulation is possible with stability of 0.2 °C. A saturated aqueous solution of zinc chloride ( $n_e = 1.51$ ) and

silicon oil ( $n_D = 1.56$ ) for low refractive OP samples as PMMA, and a saturated water solution of potassium–mercuric–iodide (KHgI) with  $n_e = 1.73$  for higher refractive PS, PC, etc. have been chosen to ensure the optical contact during the measurements. Initial estimation of the indices of refraction of the OP samples and the choice of immersion emulsions have been accomplished by means of an ellipsometric laser system LEF-3 M-1 made by Carl Zeiss Jena which measuring accuracy of  $\Delta n = \pm 0.002$  is completely insufficient to obtain precise OPs' refractometric data.

Refractive indices of OPs in VIS and NIR region are measured with the experimental set-up illustrated in Fig. 1. A G5-LOMO goniometer with an accuracy of one arc second was used with the VoF3 prism-measuring block positioned on the G5 test table. The lighting module operates with a 250 W halogen lamp applied over the entire VIS and NIR regions with interference filters (IF) made by Carl Zeiss (Jena). A new photo detector device was assembled with the aid of a plane silicon diode, operating amplifier and indicating module. The collimator forms a white light beam that falls on the fixed filter. The prism block with the OP sample is illuminated monochromatically. We found some differences in the spectral bandwidths and maxima of the filters. Therefore, the amplitude transmittances of the interference filters have been measured with the aid of Varian Carry 5 VIS-NIR spectrophotometer and all deviations were considered in presentation of the refractometric results.

The right-hand collimator with the attached photo detector determines the measuring angle  $\alpha$  (see Fig. 1). The angle of deviation  $\gamma$  is formed by the OP sample located into the  $V$ -shaped prism. The refractive index  $n_\lambda$  of the examined OP is calculated as follows:

$$n_\lambda^2 = N_\lambda^2 - \cos \gamma \sqrt{N_\lambda^2 - \cos^2 \gamma}, \quad \gamma = 90^\circ - \alpha, \quad (3.1)$$

where  $N_\lambda$  is the refractive index of the VoF3 prism,  $\gamma$  is the calculated angle of the deviated beam, and  $\alpha$  is the measured angle on the G5 set-up. The index  $N_\lambda$  of the SF3 glass is determined by the data published in [8] at standard spectral wavelengths. A new OptiColor program involving

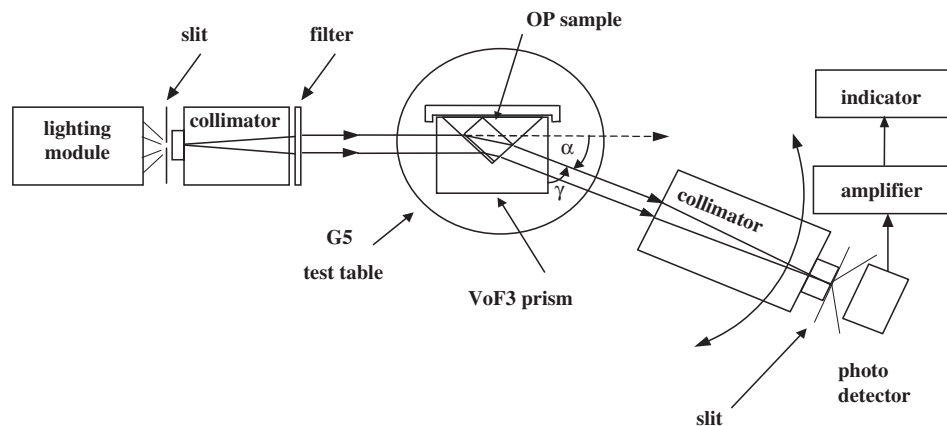


Fig. 1. Experimental set-up for measurement of the OP's refractive indices.

Table 1  
Measured refractive indices of sixteen OP materials

Material	WLs (nm)					
	435.8	486.1	587.6	703	833	1052
PMMA	1.502	1.497	1.491	1.486	1.484	1.481
PS	1.617	1.606	1.592	1.582	1.577	1.572
PC	1.612	1.599	1.585	1.575	1.570	1.565
SAN	1.588	1.578	1.567	1.558	1.554	1.550
CTE Richardson®	1.602	1.593	1.580	1.571	1.566	1.562
NAS-21®	1.593	1.584	1.571	1.564	1.558	1.554
S (low styrene)®	1.532	1.526	1.518	1.512	1.509	1.506
Optorez 1330®	1.522	1.516	1.509	1.505	1.503	1.498
Zeonex E48R®	1.543	1.538	1.531	1.526	1.523	1.520
Bayer®	1.612	1.600	1.586	1.577	1.571	1.566
Cellulose <sup>a</sup>	1.480	1.477	1.471	1.466	1.463	1.461
Polyacrylate <sup>a</sup>	1.507	1.500	1.494	1.491	1.489	1.486
Styrene <sup>a</sup>	1.534	1.527	1.519	1.513	1.510	1.507
Polycarbonate <sup>a</sup>	1.597	1.587	1.572	1.565	1.560	1.555
Polystyrene <sup>a</sup>	1.615	1.604	1.592	1.582	1.576	1.572
Acrylic <sup>a</sup>	1.502	1.498	1.492	1.488	1.485	1.483

<sup>a</sup> Control samples.

Cauchy’s dispersion formula was made to determine the dispersive coefficients of SF3 glass prism and then random refractive indices  $N_\lambda$  in VIS or NIR region are calculated.

The obtained measured data of the refractive indices of OPs at six wavelengths are presented in Table 1. The last six materials from Cellulose to Acrylic refer to laboratory specimens and are used as control samples. The presented values of measured refractive indices in Table 1 were obtained by averaging over all measured data of each series of at least five samples at given wavelength.

#### 4. Computer modelling of dispersion

In spectral regions where materials are transparent and normal dispersion occurs ( $\lambda \gg \lambda_{0i}$ ) Eq. (2.4) is reduced to Cauchy’s equation. In our previous works [7,12,25] we have applied a modified Cauchy’s approximation in the form:

$$n_\lambda^2 = A_1 + A_2\lambda^2 + \frac{A_3}{\lambda^2} + \frac{A_4}{\lambda^4} + \dots, \tag{4.1}$$

where  $A_1, A_2, A_3, A_4, \dots$  are the calculated dispersion coefficients and  $\lambda$  is the wavelength expressed in microns.

We have studied the precision of this approximation in respect to the number of the involved dispersion coefficients. They were computed with the aid of linear systems consisting of four, five, six, seven and eight equations. The available optical catalogues do not provide sufficient data for refractive indices of OPs. Because of that, we have studied the accuracy of approximation (4.1) using the extensive refractometric information on optical glasses published in Glass Catalogues. Our calculations were made on the examples of catalogue data of SF3 [8] and SF6 glass [26]. The refractive indices are given with precision of  $1 \times 10^{-5}$ . Table 2 presents the results of our calculations for the SF6 type glass.

In column 2 the catalogue refractive indices at standard spectral lines are included. Columns 3, 4, 5, 6 and 7 present the deviation  $\Delta N_\lambda$  between calculated and catalogue data. Blank places in these columns indicate the refractive indices used in the corresponding linear equation system involving 4, 5, 6, 7 or 8 dispersion coefficients in Eq. (4.1), respectively. A maximal error  $\Delta N_\lambda$  of 0.00064 and 0.00026 in the blue area of the spectrum are obtained applying four and five dispersion coefficients. We found that the accuracy to the fifth decimal place, as in SCHOTT catalogue applying the Sellmeier’s approximation, is not achievable. Using six dispersion coefficients a maximal deviation up to 0.00007 is obtained (column 5). One can see that utilization of more than six terms in the approximating row does not change significantly this result (columns 6 and 7). Therefore, we found that the usage of six dispersion coefficients in Cauchy’s approximation is sufficient to provide the accuracy of  $\pm 0.001$  of calculated refractive indices. Better precision is achievable using Sellmeier’s approximation, which should be applied in case of experimental data presented to the fourth decimal point.

Table 2  
Deviation of computed indices in respect to the catalogue data of SF6 glass

WLs (nm)	SF6 (SCHOTT)	Calculated error $\Delta N_\lambda$				
		4 coeff.	5 coeff.	6 coeff.	7 coeff.	8 coeff.
404.7 ( <i>h</i> )	1.86436	0.00064	0.00026	−0.00005	−0.00023	−0.00020
435.8 ( <i>g</i> )	1.84707					
480 ( <i>F'</i> )	1.82970	−0.00013	−0.00003	0.00001	0.00001	0.00001
486.1 ( <i>F</i> )	1.82775	−0.00014	−0.00004			
546.1 ( <i>e</i> )	1.81265	−0.00004	−0.00001	0	0	
587.6 ( <i>d</i> )	1.80518					
589.3 ( <i>D</i> )	1.80491	0	0	0	0	0
632.8 (laser)	1.79884	0.00002	0	0	0	0
643.9 ( <i>C'</i> )	1.79750	0	−0.00001	−0.00001	−0.00001	−0.00001
656.3 ( <i>C</i> )	1.79609	0.00001				
706.5 ( <i>r</i> )	1.79117					
852.1 ( <i>s</i> )	1.78157	−0.00004	0	−0.00001		
1014 ( <i>t</i> )	1.77517					
1064 (laser)	1.77380	0.00010	0.00006	0.00007	0.00006	0.00006



According to the analysis of the approximation precision best results are obtained using six dispersion coefficients in Cauchy's formula (4.1). We have realized the program OptiColor that allows us to compute the dispersion coefficients of any optical material in the region of normal dispersion. The input data consists of six refractive indices at selected measuring wavelengths (see Table 1). The dispersion coefficients from  $A_1$  to  $A_6$  can be calculated using the linear system consisting of six equations [7].

The obtained results are presented in Table 3. The computed dispersion coefficients vary in respect to the selected

wavelengths, but analysis shows that  $A_1$  always corresponds to  $n^2$  with accuracy to the first decimal place and  $A_2 \dots A_6$  introduce additional corrections of higher accuracy.

Substituting the obtained dispersion coefficients in Cauchy's dispersion formula (4.1), random refractive indices can be computed and compared with their measured values. We have calculated OPs' refractive indices at selected laser wavelengths in our earlier works [7,12,27]. In this paper we present results at some additional laser emission wavelengths. The calculated refractive data of various OPs, SF3 and SF6 glasses are given in Table 4. The Abbe

Table 3  
Computed dispersion coefficients of the examined OP materials

Material	Dispersion coefficients					
	$A_1$	$A_2$	$A_3$	$A_4$	$A_5$	$A_6$
PMMA	2.399964	-8.308636E-2	-1.919569E-1	8.720608E-2	-1.666411E-2	1.169519E-3
PS	2.610025	-6.143673E-2	-1.312267E-1	6.865432E-2	-1.295968E-2	9.055861E-4
PC	2.633127	-7.937823E-2	-1.734506E-1	8.609268E-2	-1.617892E-2	1.128933E-3
SAN	2.595568	-6.848245E-2	-1.459074E-1	7.329172E-2	-1.372433E-2	9.426682E-4
CTE Rich. <sup>®</sup>	2.663794	-1.059116E-1	-2.492271E-1	1.165541E-1	-2.211611E-2	1.545711E-3
NAS-21 <sup>®</sup>	2.054612	1.374019E-1	3.200690E-1	-1.152867E-1	2.077225E-2	-1.383569E-3
S (low styrene) <sup>®</sup>	2.360004	-4.014429E-2	-8.371568E-2	4.160019E-2	-7.586052E-3	5.071533E-4
Optorez 1330 <sup>®</sup>	2.291142	-3.311944E-2	-1.630099E-2	7.265983E-3	-6.806145E-4	1.960732E-5
Zeonex E48R <sup>®</sup>	2.482396	-6.959910E-2	-1.597726E-1	7.383333E-2	-1.398485E-2	9.728455E-4
Bayer <sup>®</sup>	2.542676	-4.366727E-2	-8.196872E-2	4.718432E-2	-8.892747E-3	6.324010E-4
Cellulose <sup>a</sup>	2.139790	-6.317682E-3	-5.920813E-3	9.613514E-3	-1.967293E-3	1.363793E-4
Polyacrylate <sup>a</sup>	2.364830	-6.955268E-2	-1.356107E-1	6.053900E-2	-1.166640E-2	8.542615E-4
Styrene <sup>a</sup>	2.274658	-5.700326E-3	-7.262838E-3	1.233343E-2	-2.481307E-3	1.784805E-4
Polycarbonate <sup>a</sup>	2.496875	-5.014035E-2	-4.188992E-2	1.732175E-2	-1.240544E-3	-1.977750E-5
Polystyrene <sup>a</sup>	2.721609	-9.982812E-2	-2.518650E-1	1.269202E-1	-2.549211E-2	1.867696E-3
Acrylic <sup>a</sup>	1.866120	2.085454E-1	4.806770E-1	-1.840693E-1	3.424849E-2	-2.340796E-3

<sup>a</sup> Control samples.

Table 4  
Abbe numbers and refractive indices of two glasses and sixteen OP materials calculated for ten laser wavelengths

Optical material	Abbe numbers		Lasing medium									
			GaN	Ar	Cu	Nd:YAG	He-Ne	Ruby	Krypton	Ti:Sapphire	Nd:YAG	Nd:YAG
	$v_d$	$v_{804}$	405 nm	488 nm	510.6 nm	532 nm	632.8 nm	694.3 nm	799.3 nm	860 nm	946 nm	1064 nm
SF3	28.14	48.48	1.7879	1.7582	1.7530	1.7488	1.7347	1.7292	1.7226	1.7189	1.7166	1.7133
SF6	25.44	44.81	1.8641	1.8272	1.8209	1.8158	1.7988	1.7923	1.7845	1.7801	1.7775	1.7737
PMMA	59.2	96.9	1.516	1.497	1.496	1.495	1.489	1.487	1.484	1.484	1.483	1.481
PS	30.5	56.6	1.634	1.605	1.602	1.599	1.587	1.583	1.578	1.576	1.574	1.572
PC	29.1	54.8	1.631	1.599	1.595	1.592	1.580	1.575	1.571	1.569	1.567	1.564
SAN	35.4	66.8	1.608	1.578	1.575	1.573	1.563	1.558	1.554	1.553	1.552	1.549
CTE Rich. <sup>®</sup>	32.8	58.5	1.612	1.592	1.589	1.587	1.576	1.571	1.567	1.566	1.564	1.562
NAS-21 <sup>®</sup>	35.5	56.3	1.588	1.583	1.580	1.577	1.568	1.564	1.559	1.557	1.555	1.554
S (low Styrene) <sup>®</sup>	44.9	79.6	1.542	1.526	1.524	1.522	1.515	1.512	1.510	1.509	1.508	1.506
Optorez 1330 <sup>®</sup>	52.0	71.9	1.531	1.516	1.514	1.513	1.507	1.505	1.503	1.502	1.501	1.498
Zeonex E48R <sup>®</sup>	56.5	100.7	1.555	1.537	1.536	1.535	1.528	1.526	1.524	1.523	1.522	1.520
Bayer <sup>®</sup>	30.0	54.5	1.623	1.599	1.596	1.593	1.582	1.578	1.572	1.570	1.568	1.566
Cellulose <sup>a</sup>	54.1	84.3	1.484	1.476	1.475	1.474	1.469	1.467	1.464	1.463	1.462	1.461
Polyacrylate <sup>a</sup>	63.3	97.8	1.521	1.499	1.498	1.497	1.492	1.491	1.490	1.488	1.487	1.485
Styrene <sup>a</sup>	42.9	77.3	1.540	1.527	1.525	1.523	1.516	1.514	1.510	1.509	1.508	1.507
Polycarbonate <sup>a</sup>	28.9	56.7	1.602	1.586	1.582	1.579	1.568	1.565	1.561	1.559	1.557	1.554
Polystyrene <sup>a</sup>	32.0	55.5	1.640	1.604	1.601	1.598	1.587	1.582	1.577	1.576	1.574	1.571
Acrylic <sup>a</sup>	57.8	97.2	1.506	1.498	1.495	1.495	1.490	1.488	1.486	1.485	1.484	1.483

<sup>a</sup> Control samples.

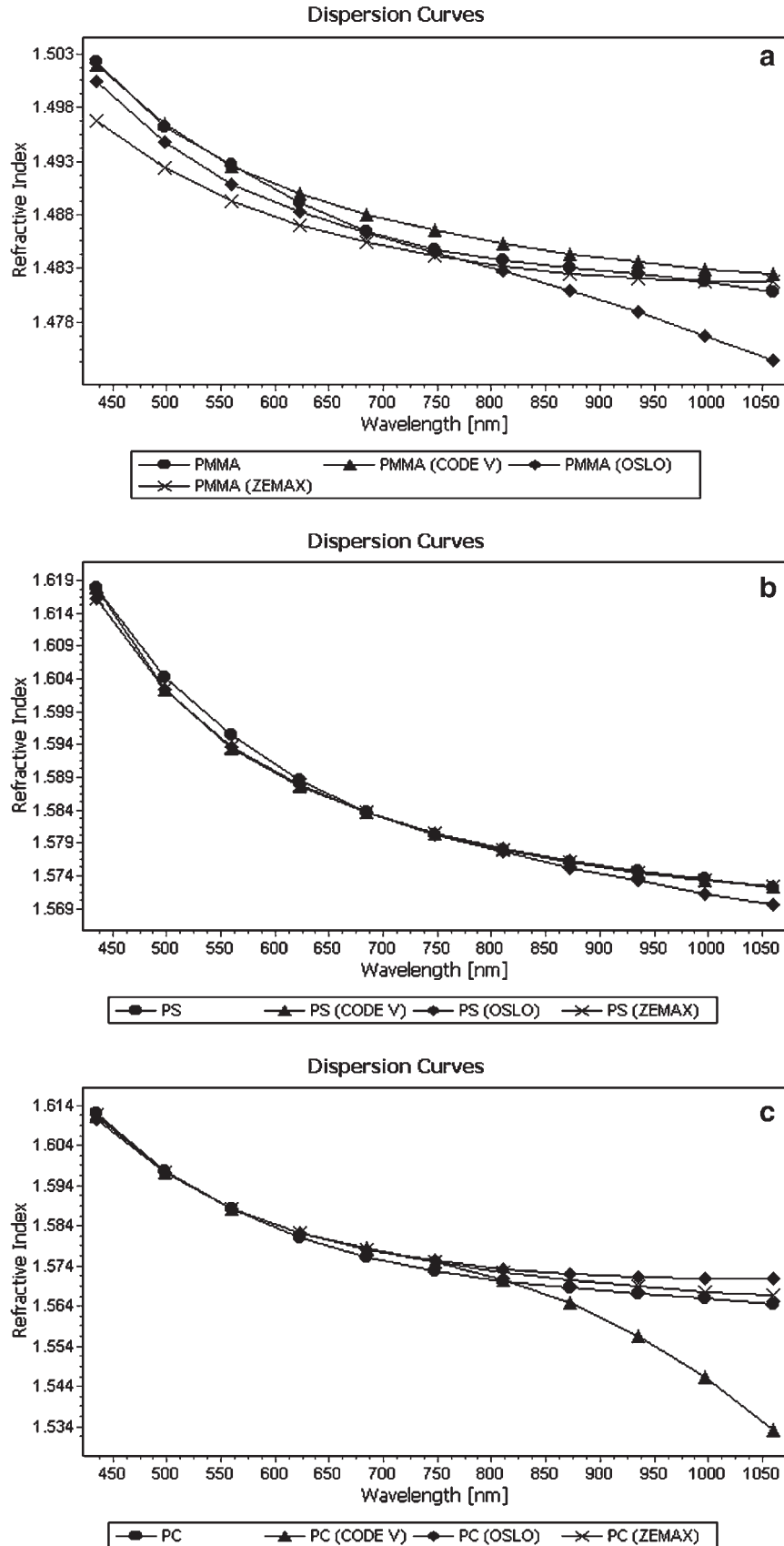


Fig. 2. Comparison of computed dispersion curves on the base of different input data for: (a) PMMA; (b) PS and (c) PC.

numbers  $v_d$  and  $v_{804}$  estimating the principal and partial dispersion, respectively, are presented in the same table.

Another option of the OptiColor program is calculation and illustration of dispersion curves of the examined material in a defined spectral area of normal dispersion. Plots of refractive indices versus wavelength for each optical polymer can be drawn. Fig. 2 presents calculated dispersion charts of the principal OPs plotted in our program on the base of our results [7] and published data of OSLO [9], ZEMAX [10] and Code V [11] program packages. A comparative analysis between the obtained dispersion curves is possible. According to approximation (4.1) the dispersion diagrams should be smooth and monotonously decreasing in the range of normal dispersion as it can be seen for dispersion curves computed on the base of our measurements. Fig. 2(a) illustrates significant differences among the obtained results in respect to the measuring method and involved approximation. Curve PMMA refers to our mea-

suring data and it is calculated by means of Eq. (4.1). It is obviously identical with curve PMMA (OSLO) in the visible part of the spectrum, but there is an intersection point between both curves in the NIR-region. This means that  $dn/d\lambda$  in both cases are different. However, our curve coincides with the PMMA – curve of ZEMAX in the NIR-spectrum. It is also obvious that curves' behaviour of measured PMMA and PMMA (CODE V) are very similar in respect to the variation of  $dn/d\lambda$ .

As it can be seen in Fig. 2(b) and (c) all calculated curves of PS and PC are strictly identical in the whole VIS region from 400 nm and further up to 800 nm. There are slight differences of PS and PC curves in the NIR (800–1100 nm) and largest deviation occurs for PC (CODE V) at 1050 nm. In general, the presented dispersion diagrams show good coincidence of the obtained results. The deviations in the NIR are probably due to the measuring accuracy of the various methods. In this spectral area the

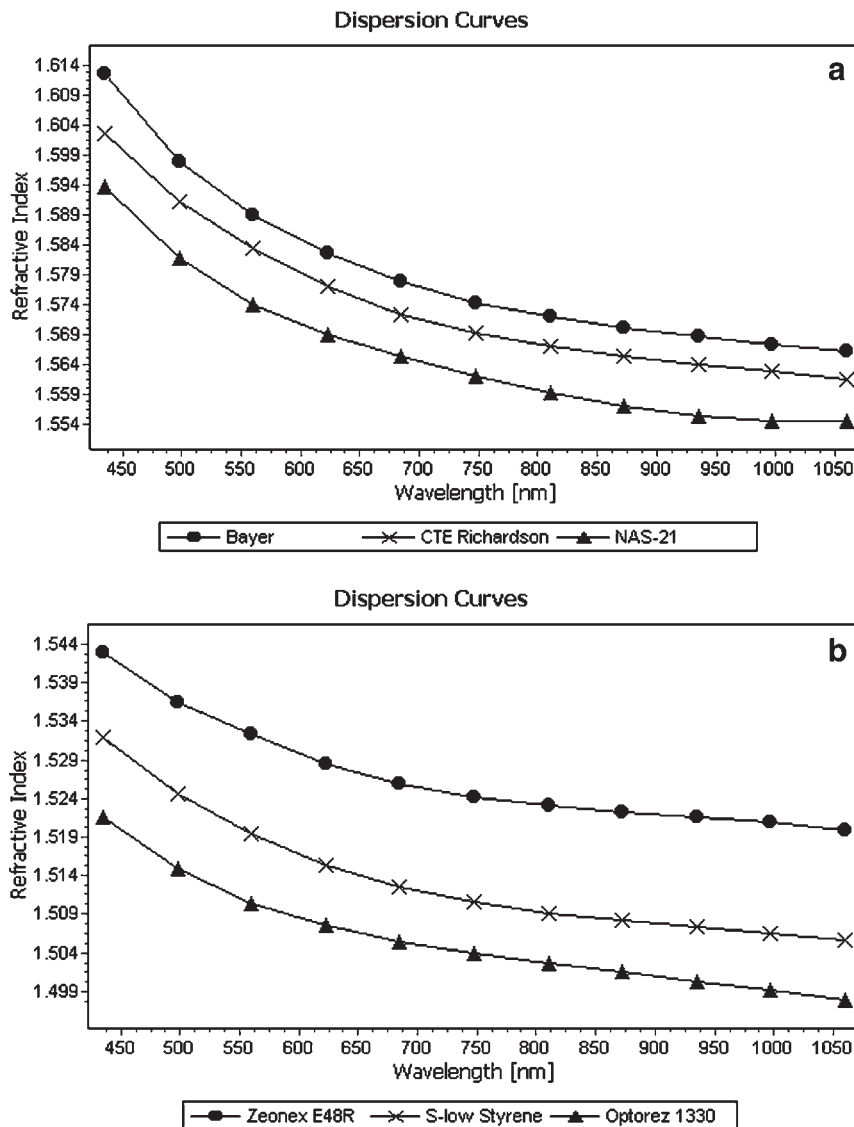


Fig. 3. Dispersion charts of selected trade-marks of OPs: (a)  $n > 1.54$  and (b)  $n < 1.54$ .

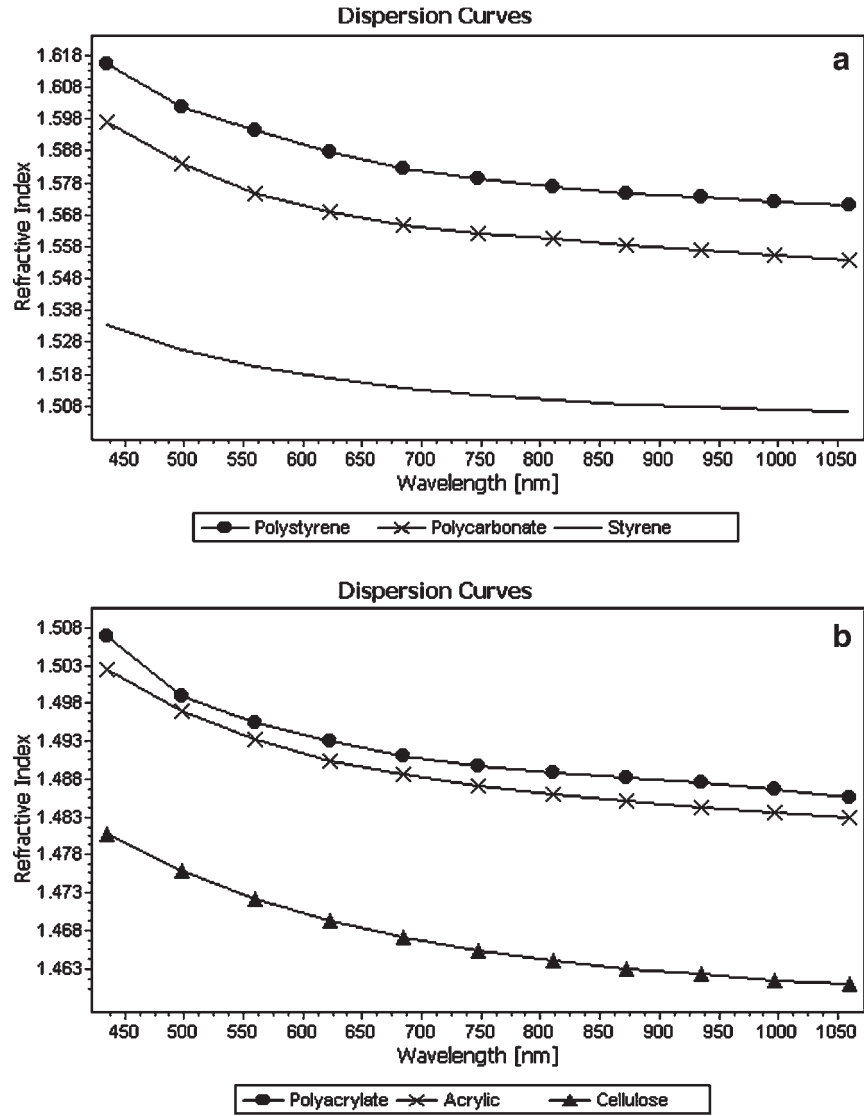


Fig. 4. Dispersion curves of control samples with: (a) higher refractive and (b) lower refractive polymers.

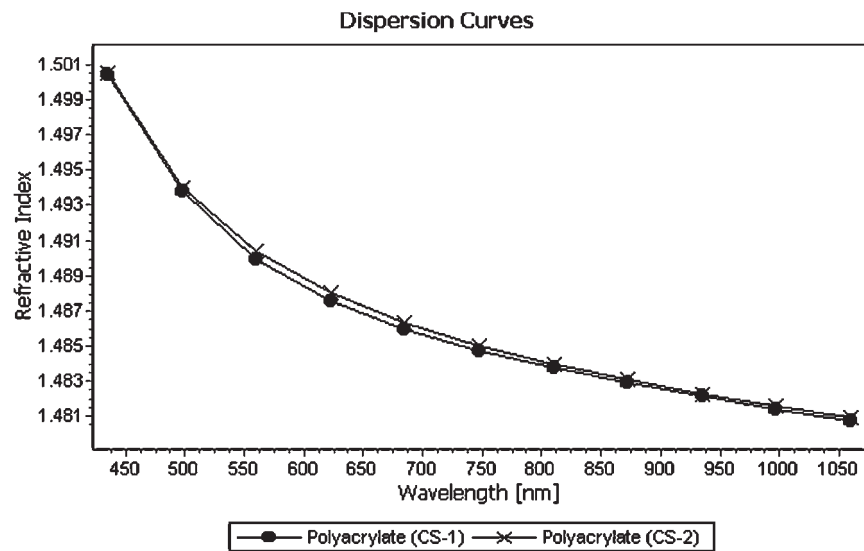


Fig. 5. The influence of birefringence and internal stresses on the refractive index values.



refractive index shows lower dependence on the wavelength and the accuracy falls down. In case of PMMA, the distinction of curves could be interpreted in a similar way. Since the refractive index is lower, stronger dependence of the measuring accuracy on the applied refractometric instruments is encountered. Obviously, birefringence and internal stresses in the samples play an additional role, too.

Computed dispersion curves on the base of measured data (Table 1) are illustrated in Figs. 3 and 4. Presentation of several dispersion curves in one diagram makes possible the comparison between the values of refractive indices and dispersion behaviour of the optical polymers. Dispersion charts of selected trade-marks (Fig. 3) and control samples (Fig. 4) are grouped into two categories: OPs with higher refractive indices (Fig. 3(a) and Fig. 4(a)) and lower refractive OP materials (Fig. 3(b) and Fig. 4(b)). The dissimilarity of the computed curves' slopes in the NIR region (see, for example, SAN, NAS-21 and Optorez 1330) is also connected with the higher measuring errors.

Internal stresses develop when a plastic part is injection moulded as a result of a temperature cycle involved [1,21]. These stresses create two indices of refraction, one in the direction of the flow and the other across the flow of the hot material. In this way birefringence arises. Bulk plastics retain residual stresses even after annealing. For example, the difference of refractive indices for PS could be as high as  $8 \times 10^{-3}$  [1,21]. The moulding process may introduce inhomogeneity too. The influence of these effects on the refractometric properties of OPs is presented for two control samples (CS-1 and CS-2) of Polyacrylate in Fig. 5. Birefringence was established in advance with the aid of a polariscope. The measuring accuracy of our refractive index measuring method and the flexibility of the new OptiColor program as well, allowed us to detect and illustrate small differences in refractive indices values.

## 5. Summary and discussions

In this paper some new measuring results for refractive and dispersion properties of OPs are presented and discussed. Refractometric measurements of sixteen American, Japanese and German OP materials were accomplished with the aid of the *V*-type prism on the Zeiss Pulfrich PR2 instrument at standard spectral lines. Additional goniometric set-up with the same *V*-type SF3 glass prism, white lighting module with interference filters, and a new sensitive photodetector device was also applied in the entire VIS and NIR spectral regions from 546 nm to 1052 nm (see Fig. 1). Our measuring results are presented with an uncertainty to the third decimal place (Table 1).

Computing accuracy of the Cauchy's approximation formula (Eq. (4.1)) is analysed on the base of glass catalogue data. Our examination shows that the usage of six dispersion coefficients ensures calculation accuracy of random refractive indices to the fourth decimal place in the region of normal dispersion (Table 2).

A new OptiColor program for computing of the dispersive properties of any optical material in the region of normal dispersion is realized. Our program calculates the dispersion coefficients (Table 3), random refractive indices and dispersion charts of the examined materials (Figs. 2–5). The Abbe numbers  $v_d$  for the principal and  $v_{804}$  for the partial dispersion are estimated. Refractive indices of the sixteen examined OPs at ten laser wavelengths in the VIS and NIR spectral areas are also calculated (Table 4).

Generation of various dispersion curves is another option of the program (Figs. 2(a)–(c)). In this way, comparative analysis of distinct input data is possible. The OptiColor program gives us the opportunity to distinguish similar dispersion curves of different samples worked out of one and the same type of material, if the input data of refractive indices differs in the third or fourth decimal place. Such variation of measured data is possible in case of inhomogeneity or birefringence of the moulded polymer samples (Fig. 5). These results demonstrate the sensitivity of the OptiColor program and prove its precision in calculations and plotting the dispersion charts.

In conclusion, the newly reported refractometric and dispersion data give detailed information about the optical properties of the examined sixteen types of OP materials. Our results could be very useful for experts from the field of physics, chemistry, optical design and photonics technology.

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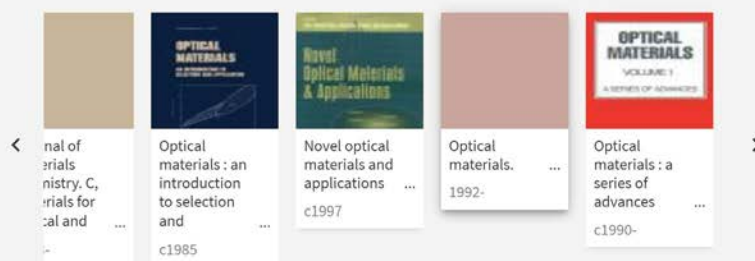
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# Noninvasive Pulse Oximetry Utilizing Skin Reflectance Photoplethysmography

YITZHAK MENDELSON, MEMBER IEEE, AND BURT D. OCHS, MEMBER IEEE

**Abstract**—The major concern in developing a sensor for reflectance pulse oximetry is the ability to measure large and stable photoplethysmograms from light which is backscattered from the skin. Utilizing a prototype optical reflectance sensor, we showed that by locally heating the skin it is possible to increase the pulsatile component of the reflected photoplethysmograms. Furthermore, we showed that additional improvements in signal-to-noise ratio can be achieved by increasing the active area of the photodetector and optimizing the separation distance between the light source and photodetector. The results from a series of *in vivo* studies to evaluate a prototype skin reflectance pulse oximeter in humans are presented.

## I. INTRODUCTION

NONINVASIVE monitoring of arterial hemoglobin oxygen saturation ( $\text{SaO}_2$ ) based upon skin reflectance spectrophotometry was first described by Brinkman and Zijlstra in 1949 [1]. They showed that changes in  $\text{SaO}_2$  can be recorded noninvasively from an optical sensor attached to the forehead. Their innovative idea to use light reflection instead of tissue transillumination, which is limited mainly to the finger tips and ear lobes, was suggested as an improvement to enable noninvasive monitoring of  $\text{SaO}_2$  from virtually any skin surface. More recent attempts to develop a skin reflectance oximeter utilizing a similar spectrophotometric approach were made by Cohen *et al.* [2] and Takatani [3]. All of those three noninvasive reflectance oximeters attempted to monitor  $\text{SaO}_2$  by measuring the absolute light intensity diffusely reflected (backscattered) from the skin.

While those developments represent significant advancements in noninvasive reflectance oximetry, limited accuracy as well as difficulties in absolute calibration were major problems with early reflectance oximeters. Although various methods have been proposed, to date, a versatile noninvasive reflectance oximeter, which can monitor  $\text{SaO}_2$  reliably from any location on the skin surface, is not yet available.

Backscattered light from living skin depends not only on the optical absorption spectrum of the blood but also on the structure and pigmentation of the skin. In an attempt to overcome this problem, Mendelson *et al.* [4]

proposed to measure  $\text{SaO}_2$  based on the principle of skin reflection photoplethysmography. We showed that  $\text{SaO}_2$  can be measured noninvasively by analyzing the pulsatile rather than the absolute, reflected light intensity  $I_r$  of the respective red and infrared photoplethysmograms according to the following empirical relationship [4]–[5]:

$$\text{SaO}_2 = A - B [I_r(\text{red})/I_r(\text{infrared})] \quad (1)$$

where  $A$  and  $B$  are empirically derived constants which are determined statistically during *in vivo* calibration in which the  $I_r(\text{red})/I_r(\text{infrared})$  ratio calculated by the pulse oximeter is compared against direct blood  $\text{SaO}_2$  measurements.  $I_r$  is obtained by a normalization process in which the pulsatile (ac) component of the red and infrared photoplethysmograms is divided by the corresponding nonpulsatile (dc) component.

In clinical applications where presently available transmission pulse oximeters cannot be used, there is a need for an optical sensor which is suitable for monitoring  $\text{SaO}_2$  utilizing light reflection from the skin. Although the principles of reflection and transmission pulse oximetry are very similar, the major limitation of reflection pulse oximetry is the comparatively low level photoplethysmograms typically recorded from the skin. The feasibility of reflection pulse oximetry, therefore, is highly dependent on the ability to detect sufficiently strong reflection photoplethysmograms.

This paper describes the considerations in designing a skin reflectance sensor for noninvasive monitoring of  $\text{SaO}_2$ . The ability to detect improved photoplethysmographic waveforms through the use of skin heating and multiple photodetectors are discussed. Results from a series of *in vivo* studies to evaluate a prototype skin reflectance pulse oximeter in humans are presented.

## II. BACKGROUND

### A. Principle of Pulse Oximetry

Pulse oximetry has been invented by Aoyagi *et al.* [6] and further refined by Nakajima *et al.* [7] and Yoshiya *et al.* [8]. This unique approach is based on the assumption that the change in light absorbed by tissue during systole is caused primarily by the arterial blood. Consequently, they showed that changes in light transmission through a pulsating vascular bed can be used to obtain an accurate noninvasive measurement of  $\text{SaO}_2$ .

The main advantage of employing a photoplethysmo-

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graphic technique is that only two wavelengths are required, thereby greatly simplifying the optical sensor. Furthermore, the requirement for blood "arterialization" which was essential in previous nonpulsatile oximeters, such as the eight wavelength Hewlett-Packard (HP) ear oximeter [9], has been eliminated. Hence, there is no need for continuous skin heating. Moreover, skin pigmentation, which can cause variable light attenuation, does not seem to affect the accuracy of pulse oximeters. This is because the ratio of the transmitted red/infrared light intensity, from which  $SaO_2$  is calculated, is obtained by a normalization process in which the ac component of the red and infrared photoplethysmograms is divided by the corresponding dc components.

The basic optical sensor of a noninvasive pulse oximeter consists of a red and infrared light emitting diodes (LED's) and a silicone photodiode. The wavelength of the red LED is typically chosen from regions of the spectra where the absorption coefficient of  $Hb$  and  $HbO_2$  are markedly different (e.g., 660 nm). The infrared wavelength, on the other hand, is typically chosen from the spectral region between 940 and 960 nm where the difference in the absorption coefficients of  $Hb$  and  $HbO_2$  is relatively small. The photodiode used has a broad spectral response that overlaps the emission spectra of the red and infrared LED's.

The light intensity detected by the photodetector depends, apart from the intensity of the incident light, mainly on the opacity of the skin, reflection by bones, tissue scattering, and the amount of blood present in the vascular bed. The amount of light attenuated by the blood varies according to the pumping action of the heart. Consequently, as tissue blood volume increases during systole, a greater portion of the incident light is absorbed by the arterial blood causing a rapidly alternating signal. Depending on the physiological state of the microvascular bed, typically, these alternating light intensity amounts to approximately 0.05–1 percent of the total light intensity either transmitted through or backscattered from the skin.

Since pulse oximeters rely on the detection of arterial pulsation, significant reduction in peripheral blood flow, such as in hypotension or hypothermia, can limit the reliability of the measurement. Nevertheless, the fact that no user calibration or site preparation is required, and the availability of small, light weight, and easy to apply sensors has made transmission pulse oximeters very popular in various clinical applications.

#### B. Reflection Versus Transmission Pulse Oximetry

In transmission pulse oximetry, sensor application is obviously limited to areas of the body, such as the finger tips, ear lobes, toes, and in infants the foot or palms where transmitted light can be readily detected. Other locations, which are not accessible to conventional transillumination techniques, i.e., the limbs, forehead, and chest may be monitored in principle using a reflection  $SaO_2$  sensor as shown schematically in Fig. 1.

Although the specific clinical utility of reflectance pulse

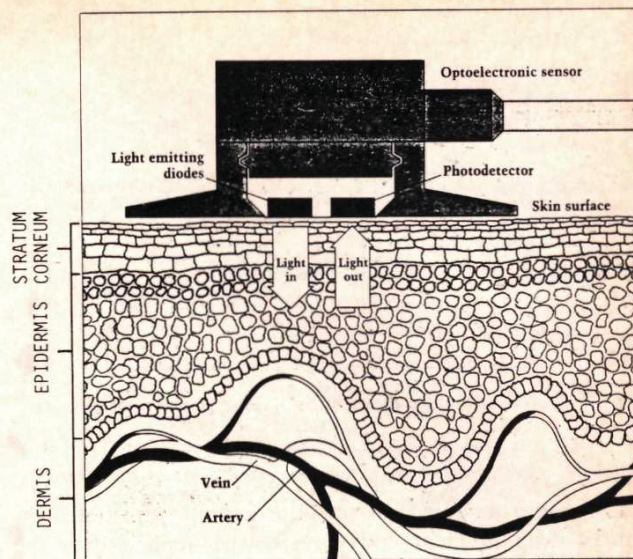


Fig. 1. Principle of reflectance pulse oximetry illustrating the optical sensor and the different layers of the skin.

oximetry has yet to be determined, it appears that the technique may have potential application for neonatal monitoring. For example, a reflectance  $SaO_2$  sensor may be of considerable value in the assessment of fetal distress during delivery if used in addition to presently available screw-type scalp ECG electrodes. Furthermore, since the skin of the chest is supplied by branches of the internal thoracic artery, which in turn stem from blood vessels leaving the aorta above the ductus arteriosus,  $SaO_2$  measurements using a reflectance sensor attached to the chest may prove to be of clinical importance when monitoring newborn infants with a patent ductus arteriosus.

### III. METHODS

#### A. Instrumentation

1) *Reflectance  $SaO_2$  Sensor:* We have constructed and tested a prototype reflectance sensor which consists of three parts: an optical sensor for monitoring  $SaO_2$ , a feedback-controlled heater for varying the local temperature of the skin under the sensor, and a laser Doppler probe for recording relative changes in skin blood flow under the sensor.

A schematic diagram illustrating the front view of the combined sensor is shown in Fig. 2. The sensor assembly can be attached to the skin by means of a double-sided, ring-shaped, tape. This attachment technique is sufficient to maintain the sensor in place without exerting excessive pressure that could significantly reduce local blood flow in the skin.

The optical sensor for monitoring  $SaO_2$  consists of red and infrared LED's with peak emission wavelength of 660 and 950 nm, respectively, and a silicone p-i-n photodiode. The half-power spectral bandwidth of each LED is approximately 20–30 nm. The LED's (dimensions:  $0.3 \times 0.3$  mm) and photodiode (dimension:  $2.0 \times 3.0$  mm)



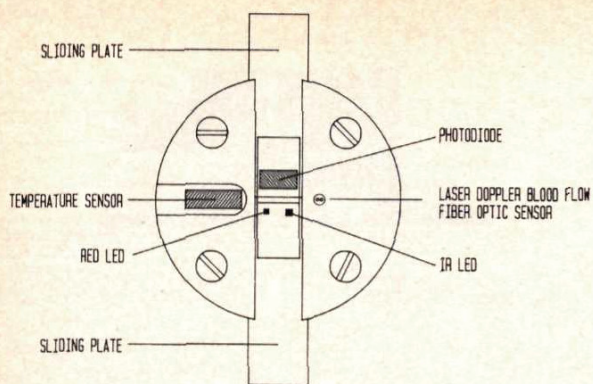


Fig. 2. Frontal view of the combined  $\text{SaO}_2$ /laser Doppler skin blood flow sensor.

chips were mounted on separate ceramic substrates. A small drop of clear epoxy resin was applied over the LED's and photodiode for protection. For investigational purposes, the ceramic substrates containing the LED's and photodiode were mounted on separate sliding plates. This arrangement provides convenient adjustment of the separation distance between the LED's and the photodiode from 4 to 11 mm. Undesired specular light reflections from the surface of the skin, as well as direct light path between the LED's and the photodiode, were minimized by recessing and optically shielding the LED's and photodiode inside the sensor assembly.

The feedback-controlled heater consists of a round thermofoil heating element (1.25 cm diameter) and a solid-state temperature transducer (Analog Devices AD590) mounted in close proximity to the surface of the sensor contacting the skin. The heater is capable of delivering a maximum power of 2 W. The temperature of the sensor can be adjusted between 34 and 45°C in 1 +/−0.1°C steps.

The distal ends of two parallel glass optical fibers (diam. 0.15 mm; separation 0.5 mm) were used for recording relative skin blood flow under the reflectance sensor. The fiber tips were mounted in close proximity to the LED's and photodiode. The proximal ends of these optical fibers were coupled to a MEDPACIFIC Model LD 5000 Laser Doppler perfusion monitor (MEDPACIFIC Corp., Seattle, WA). A 5 mW, continuous wave, HeNe laser located inside the perfusion monitor generates a monochromatic beam of red (632.8 nm) light. This light passes to the skin through one optical fiber which illuminates a region of tissue that approximates a hemisphere with a radius of about 1 mm. The light entering the tissue is scattered by the moving red blood cells causing a frequency shift proportional to the blood flow according to the Doppler principle [10]. A portion of the backscattered light from both the nonmoving tissue structures and the moving red blood cells is then collected by an adjacent optical fiber and projected onto a photodiode inside the LD 5000 monitor. The electrical output from this photodiode is processed by the perfusion monitor resulting in a continuous reading

that is proportional to the skin blood flow under the sensor. The instrument was nulled electronically before each study by adjusting the output reading to zero after the sensor was positioned over a stationary surface of white scattering material. To avoid optical interference between the LED's in the  $\text{SaO}_2$  sensor and the HeNe laser source, the reflectance pulse oximeter was turned off when skin blood flow measurements were performed.

2) *Reflectance Pulse Oximeter*: The reflectance oximeter generates digital switching pulses to drive the red and infrared LED's in the sensor alternately at a repetition rate of 1 KHz. The time multiplexed output current from the photodiode, which correspond to the red and infrared light intensities reflected from the skin, is first converted to a proportional analog voltage using a low noise operational amplifier configured as a current-to-voltage converter. The resulting output voltage is subsequently decomposed into two separate channels using two sample-and-hold circuits synchronously triggered by the same pulses driving the respective LED's. The red and infrared photoplethysmograms produced are amplified and high-pass filtered (cut-off frequency 15 Hz) to separate the ac pulses from the dc signal of each photoplethysmogram. To enable further signal processing, the respective ac and dc signals of each photoplethysmogram were digitized at a rate of 100 samples/s by an IBM-AT personal computer equipped with a Tecmar 12 bit resolution A/D-D/A data acquisition board. From the recorded signals, a computer algorithm calculates the  $I_r(\text{red})/I_r(\text{infrared})$  ratio for each heartbeat. These values are further averaged using a five-point running average algorithm. Another algorithm uses the averaged ratios to compute and display  $\text{SaO}_2$  according to (1). The  $A$  and  $B$  coefficients necessary for calculating  $\text{SaO}_2$  in the oximeter were determined previously in our laboratory based on a calibration study using the HP Model 47201A ear oximeter as a reference.

### B. In Vivo Studies

Seven Caucasian volunteers participated in the studies which were approved by our institutional review board. The subjects, five males and two females, were healthy nonsmokers ranging in age from 21 to 29 years.

To establish a reference for measuring  $\text{SaO}_2$ , we used the HP 47201A ear oximeter. The oximeter was standardized before each test by placing the ear probe in a special standardization chamber inside the ear oximeter. The ear probe was then attached to the anti-helix portion of the ear pinna with a head mount and elastic head band according to the manufacturer recommendations.

The sensor of the reflectance pulse oximeter was attached either to the volar side of the forearm or the anterior thigh region. In each case, the monitored arm or leg was immobilized in the horizontal position to minimize spurious movement artifacts.

The experimental setup used in our studies is illustrated in Fig. 3.



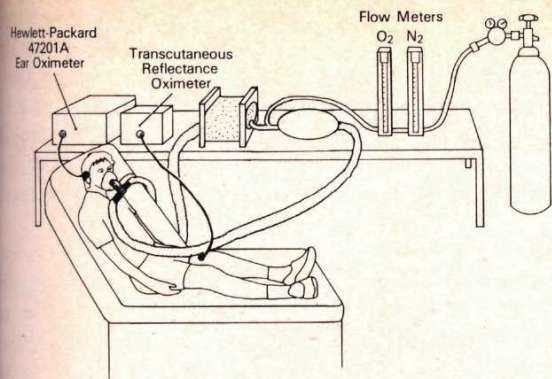


Fig. 3. Experimental setup illustrating the closed loop rebreathing circuit for obtaining different inspired O<sub>2</sub>/N<sub>2</sub> concentrations and the attachment of the oximeter sensors to the subject's ear and thigh.

IV. RESULTS

Several *in vivo* studies were performed using the prototype optical reflectance sensor and oximeter as described above. The primary objectives of the first study were to investigate the effect of 1) source/detector separation and 2) local skin heating on the pulsatile component of the red and infrared photoplethysmograms detected by the sensor. In a separate *in vivo* study, we compared SaO<sub>2</sub> values measured by the pulse oximeter from the forearm and thigh of different subjects during progressive hypoxemia with simultaneous recordings obtained from the HP ear oximeter in the range between 70-100 percent.

A. Source/Detector Separation Studies

The purpose of these studies was to determine the relationship between different LED/photodiode separations and the magnitude of the pulsatile component of each reflection photoplethysmogram. We noticed that for a constant LED intensity, the light intensity detected by the photodiode decreases roughly exponentially as the radial distance from the LED's is increased. The same basic relationship applies to both the dc and ac components of the reflected photoplethysmograms as shown in Fig. 4. This is expected since the probability that the incident photons will be absorbed as they traverse a relatively longer path length before reaching the detector is increased.

Fig. 5 shows the relative pulse amplitude of the red and infrared reflected photoplethysmograms recorded from the forearm of one subject. In this study, the incident light intensities of the red and infrared LED's were adjusted by varying the LED driving currents such that for each separation distance the dc component of each photoplethysmogram remained relatively constant. Each point represents the average values obtained for five repeated experiments performed on the same subject. In each experiment, and for each separation distance, the data acquired were averaged over a 30 s time interval.

As shown in Fig. 5, by increasing the separation distance between the LED's and photodiode from 4 to 11

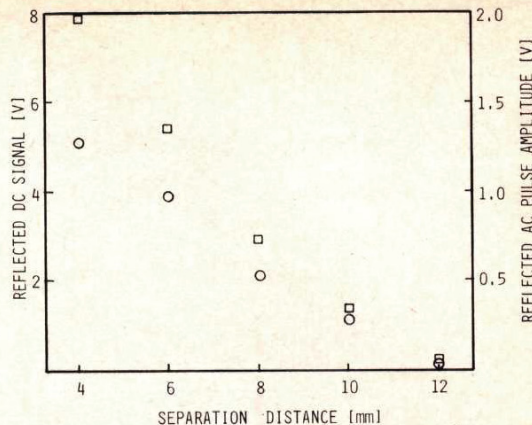


Fig. 4. The effect of LED/photodiode separation on the dc (□) and ac (○) components of the reflected infrared photoplethysmograms. Measurements were performed at a skin temperature of 43°C.

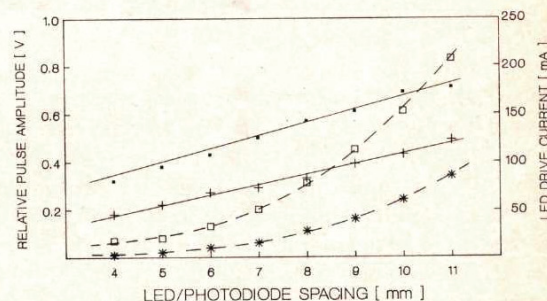


Fig. 5. Effect of LED/photodiode separation on the relative pulse amplitude of the red (+) and infrared (■) photoplethysmograms. The driving currents of the red (□) and infrared (★) LED's required to maintain a constant dc reflectance from the skin are shown for comparison.

mm, we were able to achieve almost a two-fold increase in the pulse amplitude of the infrared photoplethysmogram. Furthermore, as illustrated in Fig. 6, the mean beat-to-beat variations of the infrared photoplethysmograms, which were determined by calculating the respective coefficients of variation (i.e., the standard deviation divided by the mean for a 30 s time interval), decreased from about 7 to 3 percent. This trend indicates that the photoplethysmograms became progressively more stable as the LED/photodetector separation was increased. Similar trends were also observed for the reflected red photoplethysmograms.

B. Skin Heating Studies

Practically, it is difficult to detect large reflection photoplethysmograms from skin areas which are not very vascular, such as the chest and the limbs. In this study, we attempted to determine if local skin heating, which is known to produce vasodilatation of the microvascular bed, could be used as a practical mean to increase the pulsatile component of the reflected photoplethysmograms. Likewise, we sought to determine if skin heating could help to reduce the beat-to-beat variability in the pulsatile components of the recorded photoplethysmograms.



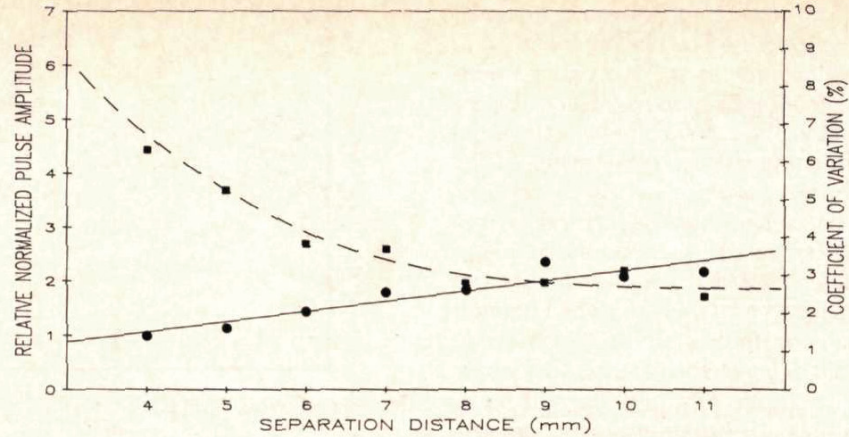


Fig. 6. Effect of LED/photodiode separation on the mean pulse amplitude (●) and the corresponding decrease in the beat-to-beat amplitude fluctuation (■) of the infrared photoplethysmograms expressed in terms of the coefficient of variation. Each pulse amplitude was normalized with respect to a separation distance of 4 mm.

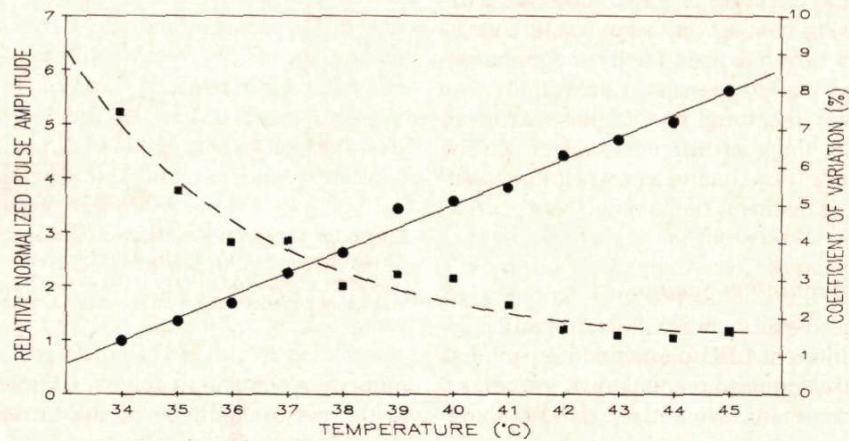


Fig. 7. Effect of skin temperature on the mean pulse amplitude (●) and the corresponding decrease in the coefficient of variation (■) of the infrared photoplethysmograms. Each pulse amplitude was normalized with respect to a separation distance of 4 mm.

Measurements were performed at a constant LED/photodiode separation of 6 mm while the subject was breathing ambient air. After attaching the reflectance sensor to the forearm, the surface of the skin was gradually heated to 45°C in 1°C step increments. The time needed to achieve a desired skin temperature depends on factors such as skin type, local blood flow, heat conductivity of the skin, and the temperature of the surrounding environment. Typically, we found that at each temperature setting, 5 min were sufficient for the skin temperature to reach steady state.

As shown in Fig. 7, by increasing the local skin temperature from 34° to 45°C, we were able to obtain a five-fold increase in the pulse amplitude of the infrared photoplethysmograms. Moreover, by heating the skin, the vascular bed under study becomes vasodilated and, there-

fore, the reflected photoplethysmograms become more stable resulting in smaller beat-to-beat amplitude fluctuations. Consequently, as our data show, the mean coefficient of variation decreased from approximately 7 to 2 percent. Similar trends were also observed for the reflected red photoplethysmograms.

The effect of local skin heating on the pulsatile component of the reflected photoplethysmograms is shown in Fig. 8. The relative skin blood flow for each temperature setting is also shown for comparison. It is clearly seen that as the temperature of the skin was increased from its initial value of 29° to 43°C, the pulse amplitude of the red and infrared photoplethysmograms increased accordingly. Furthermore, the mean pulse amplitude of the recorded waveforms remained relatively constant over a period of approximately 20 min after the heater was turned



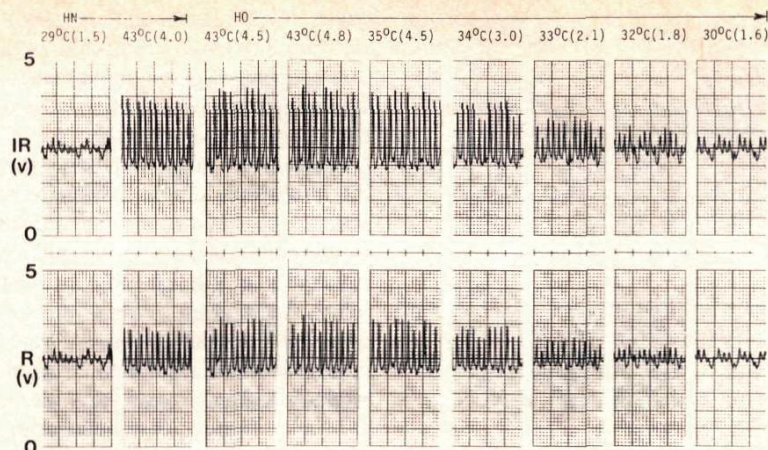


Fig. 8. Simultaneous recording of the infrared and red photoplethysmograms from the forearm at different skin temperatures. The numbers in parenthesis indicate the relative skin blood flows (scale: 0-10). Each record lasted approximately 15 s. The time elapsed between consecutive recordings is 10 min. HN = heater turned on, HO = heater turned off.

off. Thereafter, the pulse amplitude started to diminish. After about 50 min, the pulse amplitude returned to its initial level.

C. Hypoxemia Studies

Preliminary studies using our prototype reflectance sensor during progressive steady-state hypoxemia were conducted on a group of seven healthy adult volunteers.

Each subject was placed in a reclining position and asked to breathe different fractions of O<sub>2</sub>/N<sub>2</sub> gas while maintaining spontaneous respiration. The inspired O<sub>2</sub>/N<sub>2</sub> gas mixture was supplied through a fitted face mask by a closed-loop rebreathing circuit equipped with a CO<sub>2</sub> scrubber and a one-way breathing valve. The fractional inspired O<sub>2</sub> concentration (F<sub>I</sub>O<sub>2</sub>) was adjusted between 10 and 100 percent using separate gas flowmeters. The exact inspired F<sub>I</sub>O<sub>2</sub> was monitored continuously with an Instrumentation Laboratory Model 408 oxygen monitor (Instrumentation Laboratories Inc., Lexington, MA) which was inserted in the inspiratory limb.

The skin reflectance sensor was attached to the volar side of the forearm and maintained at a constant temperature of 43°C. The spacing between the LED's and the photodiode in these experiments was set to 6 mm.

Initially, the F<sub>I</sub>O<sub>2</sub> was changed in step decrements, each producing a 5 percent decrease in SaO<sub>2</sub> as measured by the reference HP ear oximeter. At each SaO<sub>2</sub> level, the inspired F<sub>I</sub>O<sub>2</sub> was maintained at a constant level until both oximeters displayed stable readings.

For each step change in F<sub>I</sub>O<sub>2</sub>, SaO<sub>2</sub> readings from our prototype reflectance pulse oximeter were averaged over 60 s time intervals and compared to the corresponding SaO<sub>2</sub> values measured simultaneously by the HP ear oximeter. The averaged readings from all seven subjects were then pooled and a linear regression analysis was performed.

A comparison between the reflectance pulse oximeter and the HP ear oximeter readings obtained from all seven subjects is shown in Fig. 9. A total of 66 pairs of data points were used in this regression analysis. Linear regression analysis of this experimental data resulted in a slope of 0.93 and a positive y intercept of 6.22 percent (r = 0.96; S.E.E = 2.20). The mean and standard deviation for the differences between the skin reflectance pulse oximeter and the HP SaO<sub>2</sub> readings were found to be -0.001 +/- 2.27, respectively.

In order to determine if repeatable SaO<sub>2</sub> measurements can be made also from body sites other than the forearm, we performed a similar series of experiments in which the sensor was applied to the thigh region of three different subjects. In these experiments, a total of 24 data points were obtained simultaneously from the reflectance pulse oximeter and the HP ear oximeter. Linear regression analysis of this data set revealed a slope of 0.93 and a positive y intercept of 6.5 percent (r = 0.99; S.E.E. = 1.56). The mean and standard deviation for the differences between the skin reflectance pulse oximeter and the HP SaO<sub>2</sub> readings were found to be -0.001 +/- 1.61, respectively.

The response of our prototype skin reflectance oximeter was further compared against simultaneous recordings of SaO<sub>2</sub> from the fingertip and earlobe made by a transmission pulse oximeter (Nellcor Model N-100, Nellcor, Inc., Hayward, CA) and the HP ear oximeter, respectively. The recordings, which are shown in Fig. 10, were obtained by asking the subject to hyperventilate and then hold his breath consecutively.

D. Multiple Photodetector Arrangement

The incident light emitted from the LED's diffuses in the skin in all directions. This is evident from the circular pattern of backscattered light surrounding the LED's. Therefore, by collecting the backscattered radiation using



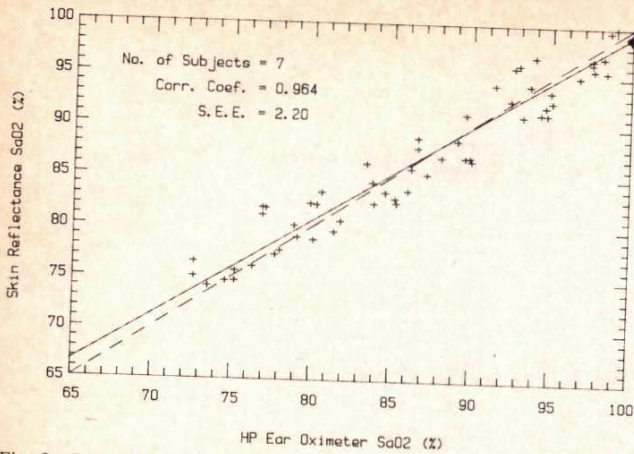


Fig. 9. Comparison of  $SaO_2$  recorded simultaneously from the forearm and the ear by the skin reflectance pulse oximeter and the HP ear oximeter, respectively.

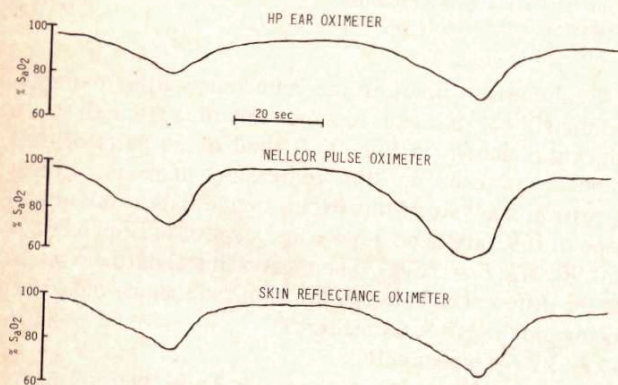


Fig. 10. Simultaneous recordings of  $SaO_2$  from the HP ear oximeter, Nellcor pulse oximeter and the prototype reflectance pulse oximeter.

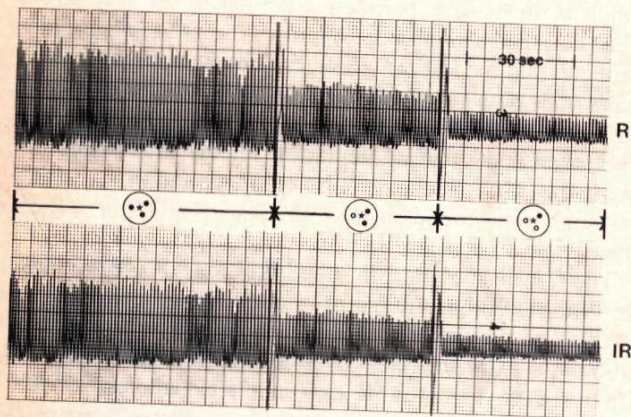


Fig. 11. Reflection photoplethysmograms recorded from the forearm using a combination of three photodiodes. The circles indicate the relative location of the photodiodes with respect to the LED's ( $\star$ ). The closed circles indicate the photodiodes which were used to collect the reflected light as shown by the corresponding traces.

several photodetectors, considerably larger photoplethysmograms could be detected.

To demonstrate the advantage of using multiple pho-

todetectors instead of only one, we modified our sensor and mounted two additional photodiodes similar in size and spectral response to that used originally. This enabled us to triple the total active area of the photodetector and thus collect a greater fraction of the backscattered light from the skin. Fig. 11 shows the spatial arrangement of all three photodiodes which were mounted symmetrically with respect to the red and infrared LED's. Also shown in this figure is the relative pulse amplitude of the red and infrared photoplethysmograms recorded from the forearm when the output currents of several photodiodes were summed simultaneously. As expected, we can see that by using multiple photodetectors a larger fraction of the backscattered radiation from the skin can be collected and, therefore, larger photoplethysmograms can be recorded.

## V. DISCUSSION

The sensor designed for this study enabled us to examine the effect of LED/photodiode separation distance as well as skin heating on the pulse amplitude of the photoplethysmograms detected by a reflectance pulse oximeter sensor.

One of the requirements in designing a reflectance pulse oximeter sensor is to determine the optimum separation distance between the LED's and the photodiode. Obviously, this distance should be selected such that photoplethysmograms with maximum pulsatile components could be detected. Generally, the pulsatile component of the reflected photoplethysmograms depends not only on the systolic blood pulse in the peripheral vascular bed but also on the amount of arterial blood within the illuminated tissue volume.

The selection of each LED driving current determines the effective penetration depth of the incident light. For a given LED/photodiode separation, it is clear that with higher levels of incident light, a larger pulsatile vascular bed will be illuminated. Consequently, the reflected photoplethysmograms will contain a larger ac component. Practical considerations, however, limit the driving current of each LED to the manufacturer specified maximum power dissipation. Alternatively, by placing the photodetector too close to the LED's, the large dc component, which is mainly due to multiple scattering of the incident photons by the blood-free stratum corneum and epidermis layers in the skin, will cause the photodetector to become saturated.

It is important to point out that although the HP ear oximeter which was used as a reference in our studies is not an acceptable primary standard for measuring  $SaO_2$ , its accuracy and reliability as a noninvasive oximeter have been widely established [11]–[13].

Our experience using the prototype reflectance sensor has shown that the pulse amplitude of the reflection photoplethysmograms depends among other factors on the position of the photodiode relative to the LED's. The selection of a particular separation distance, however, involves a tradeoff. On one hand, larger photoplethysmograms can be detected by mounting the photodiode further



apart from the LED's. On the other hand, higher LED driving currents are necessary to overcome the absorption of the incident light due to a longer optical path length.

The results of our studies also validated our hypothesis that skin heating is a feasible method for increasing the size of the reflected photoplethysmograms. We noticed that by heating the skin surface to 45°C, a five-fold increase in the pulsatile component could be achieved. We noticed also that the improvement due to skin heating can last up to 20 min from the time the temperature of the skin has reached 45°C and the heater was turned off. It is important to mention that the ability to measure accurate SaO<sub>2</sub> values by the prototype pulse oximeter sensor was independent of the exact skin temperature. We found that a minimum skin temperature of approximately 40°C is generally sufficient in order to detect adequate stable photoplethysmograms. Our experience in healthy adults has shown that at this temperature the heated sensor can remain in the same location for at least three hours without any apparent skin damage. It should be noted also that the principal objective of skin heating in our specific application is not to increase oxygen diffusion through the skin as in transcutaneous pO<sub>2</sub> monitoring although the vasodilatation effect of the applied heat on the vascular bed in the skin is basically the same.

The close correlation obtained between the prototype reflectance pulse oximeter, the HP ear oximeter and the Nellcor finger pulse oximeter is encouraging. We showed that the technique is sensitive and permits real time monitoring of SaO<sub>2</sub> from skin areas such as the forearm and the thigh. Further work, however, is needed in order to compare our reflectance pulse oximeter against SaO<sub>2</sub> measured directly from arterial blood samples and establish the potential of this technique in various clinical applications.

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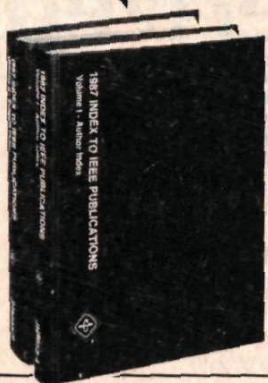


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# Noninvasive Pulse Oximetry Utilizing Skin Reflectance Photoplethysmography

YITZHAK MENDELSON, MEMBER IEEE, AND BURT D. OCHS, MEMBER IEEE

**Abstract**—The major concern in developing a sensor for reflectance pulse oximetry is the ability to measure large and stable photoplethysmograms from light which is backscattered from the skin. Utilizing a prototype optical reflectance sensor, we showed that by locally heating the skin it is possible to increase the pulsatile component of the reflected photoplethysmograms. Furthermore, we showed that additional improvements in signal-to-noise ratio can be achieved by increasing the active area of the photodetector and optimizing the separation distance between the light source and photodetector. The results from a series of *in vivo* studies to evaluate a prototype skin reflectance pulse oximeter in humans are presented.

## I. INTRODUCTION

NONINVASIVE monitoring of arterial hemoglobin oxygen saturation ( $\text{SaO}_2$ ) based upon skin reflectance spectrophotometry was first described by Brinkman and Zijlstra in 1949 [1]. They showed that changes in  $\text{SaO}_2$  can be recorded noninvasively from an optical sensor attached to the forehead. Their innovative idea to use light reflection instead of tissue transillumination, which is limited mainly to the finger tips and ear lobes, was suggested as an improvement to enable noninvasive monitoring of  $\text{SaO}_2$  from virtually any skin surface. More recent attempts to develop a skin reflectance oximeter utilizing a similar spectrophotometric approach were made by Cohen *et al.* [2] and Takatani [3]. All of those three noninvasive reflectance oximeters attempted to monitor  $\text{SaO}_2$  by measuring the absolute light intensity diffusely reflected (backscattered) from the skin.

While those developments represent significant advancements in noninvasive reflectance oximetry, limited accuracy as well as difficulties in absolute calibration were major problems with early reflectance oximeters. Although various methods have been proposed, to date, a versatile noninvasive reflectance oximeter, which can monitor  $\text{SaO}_2$  reliably from any location on the skin surface, is not yet available.

Backscattered light from living skin depends not only on the optical absorption spectrum of the blood but also on the structure and pigmentation of the skin. In an attempt to overcome this problem, Mendelson *et al.* [4]

proposed to measure  $\text{SaO}_2$  based on the principle of skin reflection photoplethysmography. We showed that  $\text{SaO}_2$  can be measured noninvasively by analyzing the pulsatile rather than the absolute, reflected light intensity  $I_r$  of the respective red and infrared photoplethysmograms according to the following empirical relationship [4]–[5]:

$$\text{SaO}_2 = A - B [I_r(\text{red})/I_r(\text{infrared})] \quad (1)$$

where  $A$  and  $B$  are empirically derived constants which are determined statistically during *in vivo* calibration in which the  $I_r(\text{red})/I_r(\text{infrared})$  ratio calculated by the pulse oximeter is compared against direct blood  $\text{SaO}_2$  measurements.  $I_r$  is obtained by a normalization process in which the pulsatile (ac) component of the red and infrared photoplethysmograms is divided by the corresponding nonpulsatile (dc) component.

In clinical applications where presently available transmission pulse oximeters cannot be used, there is a need for an optical sensor which is suitable for monitoring  $\text{SaO}_2$  utilizing light reflection from the skin. Although the principles of reflection and transmission pulse oximetry are very similar, the major limitation of reflection pulse oximetry is the comparatively low level photoplethysmograms typically recorded from the skin. The feasibility of reflection pulse oximetry, therefore, is highly dependent on the ability to detect sufficiently strong reflection photoplethysmograms.

This paper describes the considerations in designing a skin reflectance sensor for noninvasive monitoring of  $\text{SaO}_2$ . The ability to detect improved photoplethysmographic waveforms through the use of skin heating and multiple photodetectors are discussed. Results from a series of *in vivo* studies to evaluate a prototype skin reflectance pulse oximeter in humans are presented.

## II. BACKGROUND

### A. Principle of Pulse Oximetry

Pulse oximetry has been invented by Aoyagi *et al.* [6] and further refined by Nakajima *et al.* [7] and Yoshiya *et al.* [8]. This unique approach is based on the assumption that the change in light absorbed by tissue during systole is caused primarily by the arterial blood. Consequently, they showed that changes in light transmission through a pulsating vascular bed can be used to obtain an accurate noninvasive measurement of  $\text{SaO}_2$ .

The main advantage of employing a photoplethysmo-

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graphic technique is that only two wavelengths are required, thereby greatly simplifying the optical sensor. Furthermore, the requirement for blood "arterialization" which was essential in previous nonpulsatile oximeters, such as the eight wavelength Hewlett-Packard (HP) ear oximeter [9], has been eliminated. Hence, there is no need for continuous skin heating. Moreover, skin pigmentation, which can cause variable light attenuation, does not seem to affect the accuracy of pulse oximeters. This is because the ratio of the transmitted red/infrared light intensity, from which  $SaO_2$  is calculated, is obtained by a normalization process in which the ac component of the red and infrared photoplethysmograms is divided by the corresponding dc components.

The basic optical sensor of a noninvasive pulse oximeter consists of a red and infrared light emitting diodes (LED's) and a silicone photodiode. The wavelength of the red LED is typically chosen from regions of the spectra where the absorption coefficient of  $Hb$  and  $HbO_2$  are markedly different (e.g., 660 nm). The infrared wavelength, on the other hand, is typically chosen from the spectral region between 940 and 960 nm where the difference in the absorption coefficients of  $Hb$  and  $HbO_2$  is relatively small. The photodiode used has a broad spectral response that overlaps the emission spectra of the red and infrared LED's.

The light intensity detected by the photodetector depends, apart from the intensity of the incident light, mainly on the opacity of the skin, reflection by bones, tissue scattering, and the amount of blood present in the vascular bed. The amount of light attenuated by the blood varies according to the pumping action of the heart. Consequently, as tissue blood volume increases during systole, a greater portion of the incident light is absorbed by the arterial blood causing a rapidly alternating signal. Depending on the physiological state of the microvascular bed, typically, these alternating light intensity amounts to approximately 0.05–1 percent of the total light intensity either transmitted through or backscattered from the skin.

Since pulse oximeters rely on the detection of arterial pulsation, significant reduction in peripheral blood flow, such as in hypotension or hypothermia, can limit the reliability of the measurement. Nevertheless, the fact that no user calibration or site preparation is required, and the availability of small, light weight, and easy to apply sensors has made transmission pulse oximeters very popular in various clinical applications.

### B. Reflection Versus Transmission Pulse Oximetry

In transmission pulse oximetry, sensor application is obviously limited to areas of the body, such as the finger tips, ear lobes, toes, and in infants the foot or palms where transmitted light can be readily detected. Other locations, which are not accessible to conventional transillumination techniques, i.e., the limbs, forehead, and chest may be monitored in principle using a reflection  $SaO_2$  sensor as shown schematically in Fig. 1.

Although the specific clinical utility of reflectance pulse

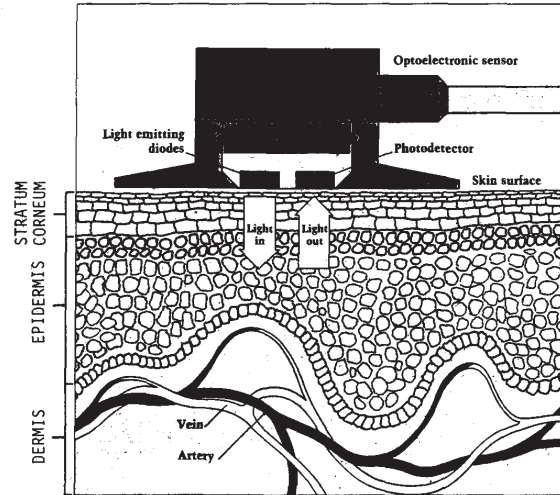


Fig. 1. Principle of reflectance pulse oximetry illustrating the optical sensor and the different layers of the skin.

oximetry has yet to be determined, it appears that the technique may have potential application for neonatal monitoring. For example, a reflectance  $SaO_2$  sensor may be of considerable value in the assessment of fetal distress during delivery if used in addition to presently available screw-type scalp ECG electrodes. Furthermore, since the skin of the chest is supplied by branches of the internal thoracic artery, which in turn stem from blood vessels leaving the aorta above the ductus arteriosus,  $SaO_2$  measurements using a reflectance sensor attached to the chest may prove to be of clinical importance when monitoring newborn infants with a patent ductus arteriosus.

## III. METHODS

### A. Instrumentation

1) *Reflectance  $SaO_2$  Sensor:* We have constructed and tested a prototype reflectance sensor which consists of three parts: an optical sensor for monitoring  $SaO_2$ , a feedback-controlled heater for varying the local temperature of the skin under the sensor, and a laser Doppler probe for recording relative changes in skin blood flow under the sensor.

A schematic diagram illustrating the front view of the combined sensor is shown in Fig. 2. The sensor assembly can be attached to the skin by means of a double-sided, ring-shaped, tape. This attachment technique is sufficient to maintain the sensor in place without exerting excessive pressure that could significantly reduce local blood flow in the skin.

The optical sensor for monitoring  $SaO_2$  consists of red and infrared LED's with peak emission wavelength of 660 and 950 nm, respectively, and a silicone p-i-n photodiode. The half-power spectral bandwidth of each LED is approximately 20–30 nm. The LED's (dimensions:  $0.3 \times 0.3$  mm) and photodiode (dimension:  $2.0 \times 3.0$  mm)

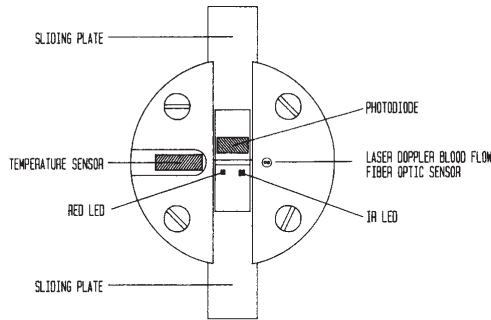


Fig. 2. Frontal view of the combined  $\text{SaO}_2$ /laser Doppler skin blood flow sensor.

chips were mounted on separate ceramic substrates. A small drop of clear epoxy resin was applied over the LED's and photodiode for protection. For investigational purposes, the ceramic substrates containing the LED's and photodiode were mounted on separate sliding plates. This arrangement provides convenient adjustment of the separation distance between the LED's and the photodiode from 4 to 11 mm. Undesired specular light reflections from the surface of the skin, as well as direct light path between the LED's and the photodiode, were minimized by recessing and optically shielding the LED's and photodiode inside the sensor assembly.

The feedback-controlled heater consists of a round thermofoil heating element (1.25 cm diameter) and a solid-state temperature transducer (Analog Devices AD590) mounted in close proximity to the surface of the sensor contacting the skin. The heater is capable of delivering a maximum power of 2 W. The temperature of the sensor can be adjusted between 34 and 45°C in 1 +/−0.1°C steps.

The distal ends of two parallel glass optical fibers (diam. 0.15 mm; separation 0.5 mm) were used for recording relative skin blood flow under the reflectance sensor. The fiber tips were mounted in close proximity to the LED's and photodiode. The proximal ends of these optical fibers were coupled to a MEDPACIFIC Model LD 5000 Laser Doppler perfusion monitor (MEDPACIFIC Corp., Seattle, WA). A 5 mW, continuous wave, HeNe laser located inside the perfusion monitor generates a monochromatic beam of red (632.8 nm) light. This light passes to the skin through one optical fiber which illuminates a region of tissue that approximates a hemisphere with a radius of about 1 mm. The light entering the tissue is scattered by the moving red blood cells causing a frequency shift proportional to the blood flow according to the Doppler principle [10]. A portion of the backscattered light from both the nonmoving tissue structures and the moving red blood cells is then collected by an adjacent optical fiber and projected onto a photodiode inside the LD 5000 monitor. The electrical output from this photodiode is processed by the perfusion monitor resulting in a continuous reading

that is proportional to the skin blood flow under the sensor. The instrument was nulled electronically before each study by adjusting the output reading to zero after the sensor was positioned over a stationary surface of white scattering material. To avoid optical interference between the LED's in the  $\text{SaO}_2$  sensor and the HeNe laser source, the reflectance pulse oximeter was turned off when skin blood flow measurements were performed.

2) *Reflectance Pulse Oximeter*: The reflectance oximeter generates digital switching pulses to drive the red and infrared LED's in the sensor alternately at a repetition rate of 1 KHz. The time multiplexed output current from the photodiode, which correspond to the red and infrared light intensities reflected from the skin, is first converted to a proportional analog voltage using a low noise operational amplifier configured as a current-to-voltage converter. The resulting output voltage is subsequently decomposed into two separate channels using two sample-and-hold circuits synchronously triggered by the same pulses driving the respective LED's. The red and infrared photoplethysmograms produced are amplified and high-pass filtered (cut-off frequency 15 Hz) to separate the ac pulses from the dc signal of each photoplethysmogram. To enable further signal processing, the respective ac and dc signals of each photoplethysmogram were digitized at a rate of 100 samples/s by an IBM-AT personal computer equipped with a Tecmar 12 bit resolution A/D-D/A data acquisition board. From the recorded signals, a computer algorithm calculates the  $I_r(\text{red})/I_r(\text{infrared})$  ratio for each heartbeat. These values are further averaged using a five-point running average algorithm. Another algorithm uses the averaged ratios to compute and display  $\text{SaO}_2$  according to (1). The *A* and *B* coefficients necessary for calculating  $\text{SaO}_2$  in the oximeter were determined previously in our laboratory based on a calibration study using the HP Model 47201A ear oximeter as a reference.

### B. In Vivo Studies

Seven Caucasian volunteers participated in the studies which were approved by our institutional review board. The subjects, five males and two females, were healthy nonsmokers ranging in age from 21 to 29 years.

To establish a reference for measuring  $\text{SaO}_2$ , we used the HP 47201A ear oximeter. The oximeter was standardized before each test by placing the ear probe in a special standardization chamber inside the ear oximeter. The ear probe was then attached to the anti-helix portion of the ear pinna with a head mount and elastic head band according to the manufacturer recommendations.

The sensor of the reflectance pulse oximeter was attached either to the volar side of the forearm or the anterior thigh region. In each case, the monitored arm or leg was immobilized in the horizontal position to minimize spurious movement artifacts.

The experimental setup used in our studies is illustrated in Fig. 3.

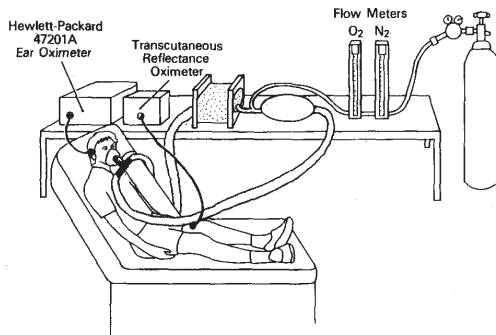


Fig. 3. Experimental setup illustrating the closed loop rebreathing circuit for obtaining different inspired  $O_2/N_2$  concentrations and the attachment of the oximeter sensors to the subject's ear and thigh.

IV. RESULTS

Several *in vivo* studies were performed using the prototype optical reflectance sensor and oximeter as described above. The primary objectives of the first study were to investigate the effect of 1) source/detector separation and 2) local skin heating on the pulsatile component of the red and infrared photoplethysmograms detected by the sensor. In a separate *in vivo* study, we compared  $SaO_2$  values measured by the pulse oximeter from the forearm and thigh of different subjects during progressive hypoxemia with simultaneous recordings obtained from the HP ear oximeter in the range between 70-100 percent.

A. Source/Detector Separation Studies

The purpose of these studies was to determine the relationship between different LED/photodiode separations and the magnitude of the pulsatile component of each reflection photoplethysmogram. We noticed that for a constant LED intensity, the light intensity detected by the photodiode decreases roughly exponentially as the radial distance from the LED's is increased. The same basic relationship applies to both the dc and ac components of the reflected photoplethysmograms as shown in Fig. 4. This is expected since the probability that the incident photons will be absorbed as they traverse a relatively longer path length before reaching the detector is increased.

Fig. 5 shows the relative pulse amplitude of the red and infrared reflected photoplethysmograms recorded from the forearm of one subject. In this study, the incident light intensities of the red and infrared LED's were adjusted by varying the LED driving currents such that for each separation distance the dc component of each photoplethysmogram remained relatively constant. Each point represents the average values obtained for five repeated experiments performed on the same subject. In each experiment, and for each separation distance, the data acquired were averaged over a 30 s time interval.

As shown in Fig. 5, by increasing the separation distance between the LED's and photodiode from 4 to 11

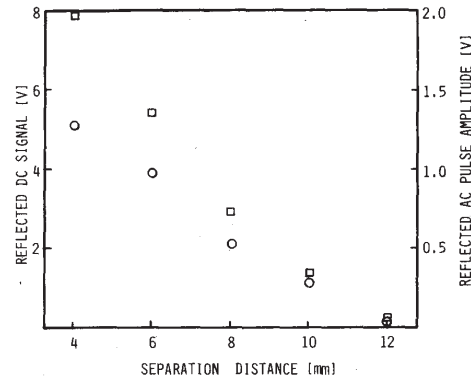


Fig. 4. The effect of LED/photodiode separation on the dc (□) and ac (○) components of the reflected infrared photoplethysmograms. Measurements were performed at a skin temperature of 43°C.

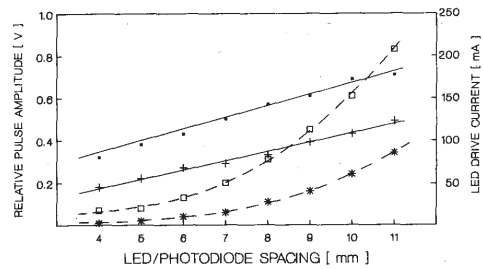


Fig. 5. Effect of LED/photodiode separation on the relative pulse amplitude of the red (+) and infrared (■) photoplethysmograms. The driving currents of the red (□) and infrared (★) LED's required to maintain a constant dc reflectance from the skin are shown for comparison.

mm, we were able to achieve almost a two-fold increase in the pulse amplitude of the infrared photoplethysmogram. Furthermore, as illustrated in Fig. 6, the mean beat-to-beat variations of the infrared photoplethysmograms, which were determined by calculating the respective coefficients of variation (i.e., the standard deviation divided by the mean for a 30 s time interval), decreased from about 7 to 3 percent. This trend indicates that the photoplethysmograms became progressively more stable as the LED/photodetector separation was increased. Similar trends were also observed for the reflected red photoplethysmograms.

B. Skin Heating Studies

Practically, it is difficult to detect large reflection photoplethysmograms from skin areas which are not very vascular, such as the chest and the limbs. In this study, we attempted to determine if local skin heating, which is known to produce vasodilatation of the microvascular bed, could be used as a practical mean to increase the pulsatile component of the reflected photoplethysmograms. Likewise, we sought to determine if skin heating could help to reduce the beat-to-beat variability in the pulsatile components of the recorded photoplethysmograms.

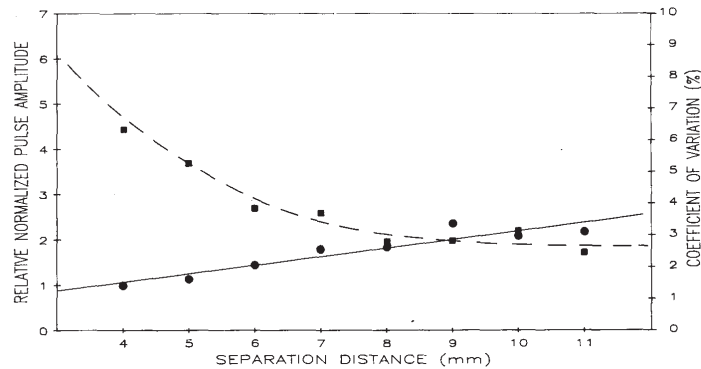


Fig. 6. Effect of LED/photodiode separation on the mean pulse amplitude (●) and the corresponding decrease in the beat-to-beat amplitude fluctuation (■) of the infrared photoplethysmograms expressed in terms of the coefficient of variation. Each pulse amplitude was normalized with respect to a separation distance of 4 mm.

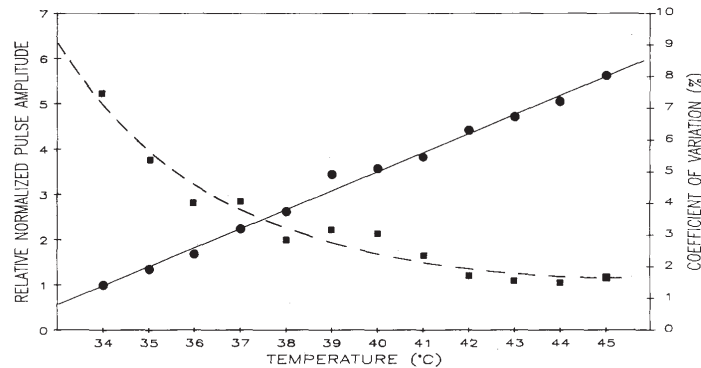


Fig. 7. Effect of skin temperature on the mean pulse amplitude (●) and the corresponding decrease in the coefficient of variation (■) of the infrared photoplethysmograms. Each pulse amplitude was normalized with respect to a separation distance of 4 mm.

Measurements were performed at a constant LED/photodiode separation of 6 mm while the subject was breathing ambient air. After attaching the reflectance sensor to the forearm, the surface of the skin was gradually heated to 45°C in 1°C step increments. The time needed to achieve a desired skin temperature depends on factors such as skin type, local blood flow, heat conductivity of the skin, and the temperature of the surrounding environment. Typically, we found that at each temperature setting, 5 min were sufficient for the skin temperature to reach steady state.

As shown in Fig. 7, by increasing the local skin temperature from 34° to 45°C, we were able to obtain a five-fold increase in the pulse amplitude of the infrared photoplethysmograms. Moreover, by heating the skin, the vascular bed under study becomes vasodilated and, there-

fore, the reflected photoplethysmograms become more stable resulting in smaller beat-to-beat amplitude fluctuations. Consequently, as our data show, the mean coefficient of variation decreased from approximately 7 to 2 percent. Similar trends were also observed for the reflected red photoplethysmograms.

The effect of local skin heating on the pulsatile component of the reflected photoplethysmograms is shown in Fig. 8. The relative skin blood flow for each temperature setting is also shown for comparison. It is clearly seen that as the temperature of the skin was increased from its initial value of 29° to 43°C, the pulse amplitude of the red and infrared photoplethysmograms increased accordingly. Furthermore, the mean pulse amplitude of the recorded waveforms remained relatively constant over a period of approximately 20 min after the heater was turned



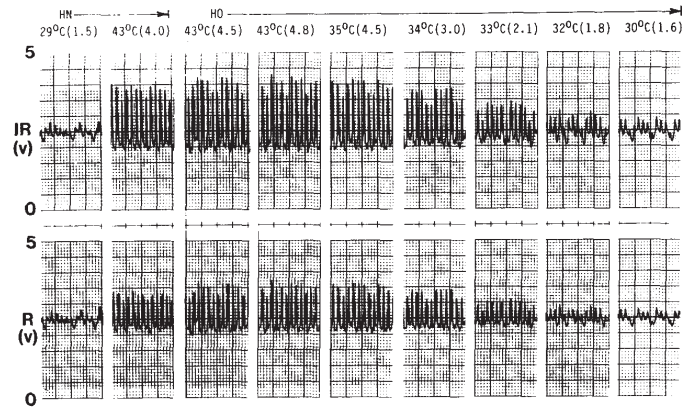


Fig. 8. Simultaneous recording of the infrared and red photoplethysmograms from the forearm at different skin temperatures. The numbers in parenthesis indicate the relative skin blood flows (scale: 0-10). Each record lasted approximately 15 s. The time elapsed between consecutive recordings is 10 min. *HN* = heater turned on, *HO* = heater turned off.

off. Thereafter, the pulse amplitude started to diminish. After about 50 min, the pulse amplitude returned to its initial level.

### C. Hypoxemia Studies

Preliminary studies using our prototype reflectance sensor during progressive steady-state hypoxemia were conducted on a group of seven healthy adult volunteers.

Each subject was placed in a reclining position and asked to breathe different fractions of  $O_2/N_2$  gas while maintaining spontaneous respiration. The inspired  $O_2/N_2$  gas mixture was supplied through a fitted face mask by a closed-loop rebreathing circuit equipped with a  $CO_2$  scrubber and a one-way breathing valve. The fractional inspired  $O_2$  concentration ( $F_I O_2$ ) was adjusted between 10 and 100 percent using separate gas flowmeters. The exact inspired  $F_I O_2$  was monitored continuously with an Instrumentation Laboratory Model 408 oxygen monitor (Instrumentation Laboratories Inc., Lexington, MA) which was inserted in the inspiratory limb.

The skin reflectance sensor was attached to the volar side of the forearm and maintained at a constant temperature of  $43^\circ C$ . The spacing between the LED's and the photodiode in these experiments was set to 6 mm.

Initially, the  $F_I O_2$  was changed in step decrements, each producing a 5 percent decrease in  $SaO_2$  as measured by the reference HP ear oximeter. At each  $SaO_2$  level, the inspired  $F_I O_2$  was maintained at a constant level until both oximeters displayed stable readings.

For each step change in  $F_I O_2$ ,  $SaO_2$  readings from our prototype reflectance pulse oximeter were averaged over 60 s time intervals and compared to the corresponding  $SaO_2$  values measured simultaneously by the HP ear oximeter. The averaged readings from all seven subjects were then pooled and a linear regression analysis was performed.

A comparison between the reflectance pulse oximeter and the HP ear oximeter readings obtained from all seven subjects is shown in Fig. 9. A total of 66 pairs of data points were used in this regression analysis. Linear regression analysis of this experimental data resulted in a slope of 0.93 and a positive y intercept of 6.22 percent ( $r = 0.96$ ;  $S.E.E. = 2.20$ ). The mean and standard deviation for the differences between the skin reflectance pulse oximeter and the HP  $SaO_2$  readings were found to be  $-0.001 \pm 2.27$ , respectively.

In order to determine if repeatable  $SaO_2$  measurements can be made also from body sites other than the forearm, we performed a similar series of experiments in which the sensor was applied to the thigh region of three different subjects. In these experiments, a total of 24 data points were obtained simultaneously from the reflectance pulse oximeter and the HP ear oximeter. Linear regression analysis of this data set revealed a slope of 0.93 and a positive y intercept of 6.5 percent ( $r = 0.99$ ;  $S.E.E. = 1.56$ ). The mean and standard deviation for the differences between the skin reflectance pulse oximeter and the HP  $SaO_2$  readings were found to be  $-0.001 \pm 1.61$ , respectively.

The response of our prototype skin reflectance oximeter was further compared against simultaneous recordings of  $SaO_2$  from the fingertip and earlobe made by a transmission pulse oximeter (Nellcor Model N-100, Nellcor, Inc., Hayward, CA) and the HP ear oximeter, respectively. The recordings, which are shown in Fig. 10, were obtained by asking the subject to hyperventilate and then hold his breath consecutively.

### D. Multiple Photodetector Arrangement

The incident light emitted from the LED's diffuses in the skin in all directions. This is evident from the circular pattern of backscattered light surrounding the LED's. Therefore, by collecting the backscattered radiation using

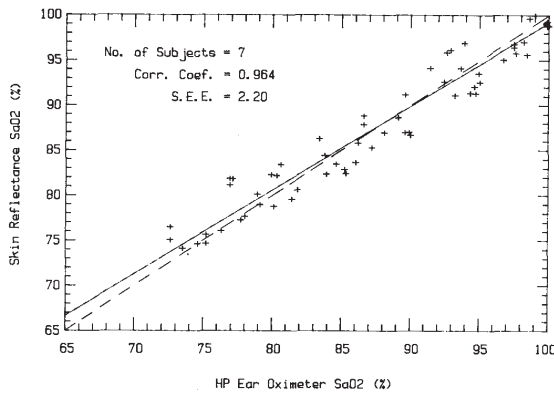


Fig. 9. Comparison of  $\text{SaO}_2$  recorded simultaneously from the forearm and the ear by the skin reflectance pulse oximeter and the HP ear oximeter, respectively.

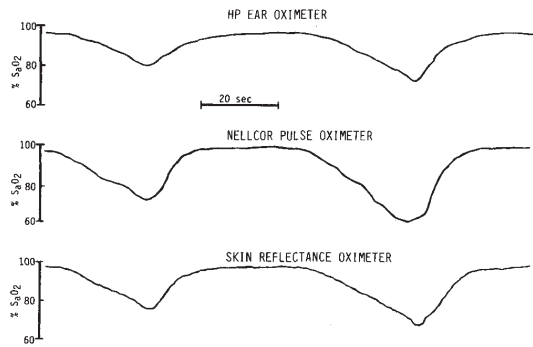


Fig. 10. Simultaneous recordings of  $\text{SaO}_2$  from the HP ear oximeter, Nellcor pulse oximeter and the prototype reflectance pulse oximeter.

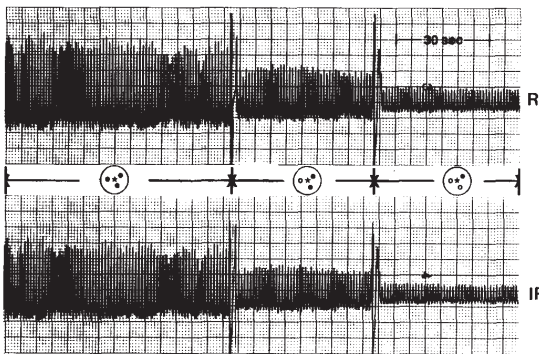


Fig. 11. Reflection photoplethysmograms recorded from the forearm using a combination of three photodiodes. The circles indicate the relative location of the photodiodes with respect to the LED's (\*). The closed circles indicate the photodiodes which were used to collect the reflected light as shown by the corresponding traces.

several photodetectors, considerably larger photoplethysmograms could be detected.

To demonstrate the advantage of using multiple pho-

todetectors instead of only one, we modified our sensor and mounted two additional photodiodes similar in size and spectral response to that used originally. This enabled us to triple the total active area of the photodetector and thus collect a greater fraction of the backscattered light from the skin. Fig. 11 shows the spatial arrangement of all three photodiodes which were mounted symmetrically with respect to the red and infrared LED's. Also shown in this figure is the relative pulse amplitude of the red and infrared photoplethysmograms recorded from the forearm when the output currents of several photodiodes were summed simultaneously. As expected, we can see that by using multiple photodetectors a larger fraction of the backscattered radiation from the skin can be collected and, therefore, larger photoplethysmograms can be recorded.

## V. DISCUSSION

The sensor designed for this study enabled us to examine the effect of LED/photodiode separation distance as well as skin heating on the pulse amplitude of the photoplethysmograms detected by a reflectance pulse oximeter sensor.

One of the requirements in designing a reflectance pulse oximeter sensor is to determine the optimum separation distance between the LED's and the photodiode. Obviously, this distance should be selected such that photoplethysmograms with maximum pulsatile components could be detected. Generally, the pulsatile component of the reflected photoplethysmograms depends not only on the systolic blood pulse in the peripheral vascular bed but also on the amount of arterial blood within the illuminated tissue volume.

The selection of each LED driving current determines the effective penetration depth of the incident light. For a given LED/photodiode separation, it is clear that with higher levels of incident light, a larger pulsatile vascular bed will be illuminated. Consequently, the reflected photoplethysmograms will contain a larger ac component. Practical considerations, however, limit the driving current of each LED to the manufacturer specified maximum power dissipation. Alternatively, by placing the photodetector too close to the LED's, the large dc component, which is mainly due to multiple scattering of the incident photons by the blood-free stratum corneum and epidermis layers in the skin, will cause the photodetector to become saturated.

It is important to point out that although the HP ear oximeter which was used as a reference in our studies is not an acceptable primary standard for measuring  $\text{SaO}_2$ , its accuracy and reliability as a noninvasive oximeter have been widely established [11]-[13].

Our experience using the prototype reflectance sensor has shown that the pulse amplitude of the reflection photoplethysmograms depends among other factors on the position of the photodiode relative to the LED's. The selection of a particular separation distance, however, involves a tradeoff. On one hand, larger photoplethysmograms can be detected by mounting the photodiode further

apart from the LED's. On the other hand, higher LED driving currents are necessary to overcome the absorption of the incident light due to a longer optical path length.

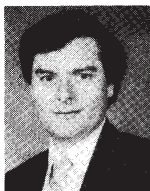
The results of our studies also validated our hypothesis that skin heating is a feasible method for increasing the size of the reflected photoplethysmograms. We noticed that by heating the skin surface to 45°C, a five-fold increase in the pulsatile component could be achieved. We noticed also that the improvement due to skin heating can last up to 20 min from the time the temperature of the skin has reached 45°C and the heater was turned off. It is important to mention that the ability to measure accurate SaO<sub>2</sub> values by the prototype pulse oximeter sensor was independent of the exact skin temperature. We found that a minimum skin temperature of approximately 40°C is generally sufficient in order to detect adequate stable photoplethysmograms. Our experience in healthy adults has shown that at this temperature the heated sensor can remain in the same location for at least three hours without any apparent skin damage. It should be noted also that the principal objective of skin heating in our specific application is not to increase oxygen diffusion through the skin as in transcutaneous pO<sub>2</sub> monitoring although the vasodilatation effect of the applied heat on the vascular bed in the skin is basically the same.

The close correlation obtained between the prototype reflectance pulse oximeter, the HP ear oximeter and the Nellcor finger pulse oximeter is encouraging. We showed that the technique is sensitive and permits real time monitoring of SaO<sub>2</sub> from skin areas such as the forearm and the thigh. Further work, however, is needed in order to compare our reflectance pulse oximeter against SaO<sub>2</sub> measured directly from arterial blood samples and establish the potential of this technique in various clinical applications.

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