
WIRELESS COMMUNICATIONS
DESIGN HANDBOOK
*Aspects of Noise, Interference,
and Environmental Concerns*

VOLUME 1: SPACE INTERFERENCE

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Chapter 6 | Noise Representations in Transponders and Multiple Access

6.0 Introduction

At the heart of a satellite communications system is the transponder. The transponder consists of input and output filters, up and down converters, phase-locked loops, and traveling wave tube amplifiers (TWTAs.) More modern transponder systems are using solid state power amplifiers (SSPAs). We now consider the nonlinear behavior of the transponder. A block diagram representation of a typical transponder was shown in Figure 5.36.

Let the input of the transponder be represented by

$$S_i(t) = A \cos(\omega_c t + f(\theta)), \quad (6.1)$$

where $\omega_c = 2\pi f_c$ is the angular carrier frequency and phase of the input signal. The transponder output can be represented as

$$S_{out} = g(A)\cos(\omega_c t + f(\theta) + f(A)). \quad (6.2)$$

If $g(A)$ and $f(A)$ are independent of ω_c , let $A_k(t)$, $\omega_c + \omega_k$, and $f_k(\theta(t))$ denote the envelope, the angular carrier frequency, and the phase of the k th carrier. ω_c is the midband frequency or center frequency, which can take any value within a transponder bandwidth. For m number of modulated carriers, access to a transponder input can be represented by

$$\begin{aligned} S_i(t) &= \sum_{k=1}^m A_k(t) \cos\{(\omega_c + \omega_k)t + f_k(\theta(t))\} = \left[\sum_{k=1}^m A_k(t) \cos(\omega_k t + f_k(\theta(t))) \right] \cos \omega_c t \\ &\quad - \left[\sum_{k=1}^m A_k(t) \sin(\omega_k t + f_k(\theta(t))) \right] \sin \omega_c t \\ &= X(t) \cos \omega_c t - Y(t) \sin \omega_c t = \sqrt{X^2 + Y^2} \cos \left[\omega_c t + \tan^{-1} \frac{Y}{X} \right] \\ &= \text{Re} \left\{ \sqrt{X^2 + Y^2} \exp \left[j\omega_c t + j \tan^{-1} \left(\frac{Y}{X} \right) \right] \right\}. \end{aligned} \quad (6.3)$$

The corresponding transponder output is

$$\begin{aligned} S_{\text{out}}(t) &= \text{Re} \left\{ g(\sqrt{X^2 + Y^2}) \cdot \exp \left[j\omega_c t + \tan^{-1} \left(\frac{Y}{X} \right) + jf(\sqrt{X^2 + Y^2}) \right] \right\} \\ &= \text{Re} \left\{ g(\sqrt{X^2 + Y^2}) \cdot \exp \left[jf(\sqrt{X^2 + Y^2}) \right] \frac{(X + jY)}{\sqrt{X^2 + Y^2}} \exp(j\omega_c t) \right\}. \end{aligned} \quad (6.4)$$

Define the double Fourier transform

$$\begin{aligned} L(u, v) &= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \frac{g(\sqrt{X^2 + Y^2})}{\sqrt{X^2 + Y^2}} \exp \left[jf(\sqrt{X^2 + Y^2}) \right] \\ &\quad (X + jY) \exp[-juX - jvY] dx dy, \end{aligned}$$

which means

$$\begin{aligned} g(\sqrt{X^2 + Y^2}) \exp \left[jf(\sqrt{X^2 + Y^2}) \right] \frac{(X + jY)}{\sqrt{X^2 + Y^2}} \\ = \frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} L(u, v) \exp[juX + jvY] du dv. \end{aligned}$$

Therefore,

$$S_{\text{out}}(t) = \frac{1}{(2\pi)^2} \text{Re} \left[\exp(j\omega_c t) \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} L(u, v) \exp[juX + jvY] du dv \right].$$

After further mathematical manipulations,

$$\begin{aligned} S_{\text{out}}(t) &= \text{Re} \left\{ \exp(j\omega_c t) \sum_{K_1=-\infty}^{\infty} \sum_{K_2=-\infty}^{\infty} \dots \sum_{K_n=-\infty}^{\infty} \right. \\ &\quad \left. \exp \left[jK_1(\omega_1 t + f(\theta_1)) + jK_2(\omega_2 t + f(\theta_2)) \right. \right. \\ &\quad \left. \left. + \dots + jK_n(\omega_n t + f(\theta_n)) \right] \cdot N(k) \right\}, \end{aligned}$$

where K_1, K_2, \dots, K_n can be zero or any integers either positive or negative, and

$$\begin{aligned} N(k) &= \frac{1}{(2\pi)^2} \int_{-\infty}^{\infty} \frac{g(\sqrt{X^2 + Y^2})}{\sqrt{X^2 + Y^2}} \exp[jf(\sqrt{X^2 + Y^2})] (X + jY) \prod_{\ell=1}^m J_{K_\ell}(A_\ell \sqrt{u^2 + v^2}) \\ &\quad \times \exp \left[j \sum_{\ell=1}^m K_\ell \tan^{-1} \frac{u}{v} \right] \exp[-jux - jvy] dx dy du dv, \end{aligned} \quad (6.5)$$

where J_K is the K th-order Bessel function. Using the polar coordinate transformation

$$\begin{aligned} X &= \rho \cos \xi, & u &= \gamma \sin \eta \\ Y &= \rho \sin \xi, & v &= \gamma \cos \eta \end{aligned}$$

and performing the integration on ξ and η simplifies the expression to the following two cases. For $K_1 + K_2 + \dots + K_n = 1$,

$$N(k) = \int_0^{\infty} \int_0^{\infty} \gamma \prod_{\ell=1}^m J_{K_\ell}(A_\ell \gamma) \rho g(\rho) \exp(jf(\rho)) \cdot J_1(\gamma \rho) d\gamma d\rho$$

and for $K_1 + K_2 + \dots + K_n \neq 1$, $N(k) = 0$.

Finally, the output of the transponder can be expressed as

$$S_{\text{out}}(t) = \text{Re}\{N(k) \exp[j(\bar{\omega} t + f_{\bar{\ell}}^{-}(\theta))]\}, \quad (6.6)$$

where

$$\begin{aligned} f_{\bar{\ell}}^{-}(\theta) &= \sum_{\ell=1}^m K_\ell f_\ell(\theta) \\ \bar{\omega} &= \omega_c + \sum_{\ell=1}^m K_\ell \omega_\ell. \end{aligned}$$

In the case of numerical computation, the factor $g(\rho)\exp[jf(\rho)]$ can be approximated by

$$g(\rho)\exp[jf(\rho)] = \sum_{\ell=1}^L b_\ell J_1(\alpha \ell \rho),$$

where L is the number of coefficients needed, J_1 is the first-order Bessel function, and $\alpha = 2\pi/(\text{period of the Fourier series})$. For a given transponder, the characteristics $g(\rho)$ and $f(\rho)$ are known. Since $g(\rho)$ and $f(\rho)$ are given, the coefficients b_ℓ can be obtained by an approximation, and $N(k)$ reduces to

$$N(k) = \sum_{\ell=1}^L b_\ell \prod_{\ell=1}^m J_{K_\ell}(\alpha \ell A_\ell), \quad (6.7)$$

which outlines the amplitude for each input signal. The b_ℓ are determined by best fit from the input data of $g(\rho)$ and $f(\rho)$ in terms of least-square error. Computer programs can calculate $S_{\text{out}}(t)$ when $S_i(t)$, $g(\rho)$, and $f(\rho)$ are given.

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